

EARLY ONLINE RELEASE

This is a PDF of a manuscript that has been peer-reviewed and accepted for publication. As the article has not yet been formatted, copy edited or proofread, the final published version may be different from the early online release.

This pre-publication manuscript may be downloaded, distributed and used under the provisions of the Creative Commons Attribution 4.0 International (CC BY 4.0) license. It may be cited using the DOI below.

The DOI for this manuscript is

DOI:10.2151/jmsj.2025-008

J-STAGE Advance published date: December 26, 2024

The final manuscript after publication will replace the preliminary version at the above DOI once it is available.

1 **Recent research and operational tools for improved**
2 **understanding and diagnosis of tropical cyclone inner core**
3 **structure**

4
5 **Kosuke Ito***

6 *Disaster Prevention Research Institute, Kyoto University, Uji, Japan*

7 *Typhoon Science and Technology Research Center, Yokohama National University,*
8 *Yokohama, Japan*

9
10 **Yoshiaki Miyamoto**

11 *Faculty of Environment and Information Studies, Keio University, Fujisawa, Japan,*

12 *RIKEN Center for Computational Science, Kobe, Japan,*

13 *Typhoon Science and Technology Research Center, Yokohama National University,*
14 *Yokohama, Japan*

15
16 **Chun-Chieh Wu**

17 *Department of Atmospheric Sciences, National Taiwan University, Taipei, Taiwan*

18
19 **Anthony Didlake**

20 *Department of Meteorology and Atmospheric Science, Pennsylvania State University, PA,*
21 *USA*

22
23 **James Hlywiak**

24 *National Research Council, Monterey, CA, USA*

25 *Naval Research Laboratory, Monterey, CA, USA*

26

27 **Yi-Hsuan Huang**

28 *Department of Atmospheric Sciences, National Taiwan University, Taipei, Taiwan*

29

30 **Tsz-Kin Lai**

31 *Space and Atmospheric Physics Group, Department of Physics, Imperial College London,*

32 *London, UK*

33 *Meteorological Research Division, Environment and Climate Change Canada, Dorval,*

34 *Quebec, Canada*

35

36 **Lauren Pattie**

37 *Environmental Prediction Services - Tropical Cyclones, Australian Bureau of Meteorology,*

38 *Brisbane, Australia*

39

40 **Nannan Qin**

41 *Nanjing Joint Institute for Atmospheric Sciences, Chinese Academy of Meteorological*

42 *Sciences, Nanjing, China*

43

44 **Udai Shimada**

45 *Meteorological Research Institute, Japan Meteorological Agency, Tsukuba, Japan*

46

47 **Dandan Tao**

48 *Geophysical Institute, University of Bergen, Bergen, Norway*

49

Yohei Yamada

50

Research Institute for Global Change, Japan Agency for Marine-Earth Science and

51

Technology, Yokohama, Japan

52

53

Jun A. Zhang

54

Hurricane Research Division, Atlantic Oceanographic and Meteorological Laboratory,

55

Miami, FL, NOAA

56

Cooperative Institute for Marine and Atmospheric Studies, University of Miami, Miami, FL,

57

USA

58

59

Sachie Kanada

60

Institute for Space-Earth Environmental Research, Nagoya University, Nagoya, Japan

61

62

Derrick Herndon

63

Cooperative Institute for Meteorological Satellite Studies, University of Wisconsin-Madison,

64

Madison, WI, USA

65

66

Revision submitted to Journal of the Meteorological Society of Japan

67

13 December 2024

68

69

70

** Corresponding author: Kosuke Ito, Kyoto University, Uji, Japan*

71

Email: ito.kosuke.2i@kyoto-u.ac.jp

72

Tel: +81-774-38-4159

73

Fax: +81-774-38-4158

75
76
77
78
79
80
81
82
83
84
85
86
87
88
89
90
91

Abstract

The inner core of a tropical cyclone (TC) is vital for TC energetics and often undergoes dramatic changes. This article provides a review on the understanding and operational practices of the structural changes in the TC inner core, mainly focusing on recent literature and activities. The inner core structure of a TC is generally described as an axisymmetric vortex in the vicinity of a hydrostatic and gradient wind-balanced state. However, this schematic can sometimes be oversimplified. Recent studies have documented small-scale features of the inner core, structural changes in TC rapid intensification, secondary eyewall formation, and eyewall replacement cycles using observational data, and idealized and sophisticated models. In line with the progress in understanding the inner core structure, several operational agencies have recently analyzed TC structural changes using their subjective analyses or diagnostic tools, contributing to disaster prevention. We also discuss potential impacts of climate change on the inner core structure, for which further work is required to reach a solid conclusion.

Keywords: Tropical cyclones; structure change; operational forecast; climate change

93 1. Introduction

94 Tropical cyclones (TCs; all acronyms are summarized in the appendix) are often
95 highly destructive, and an understanding of these systems and their operational activities
96 are important in atmospheric science and disaster prevention. The inner core of TCs is
97 particularly important owing to the strong winds and intense rainfall. The inner core of a
98 developed TC consists of an eye, an eyewall surrounding the eye, and inner rainbands
99 originating from the eyewalls. It is generally accepted that diabatic heating in the eyewall,
100 banking upward and outward, is vital for TC energetics, and the relevant updraft forces
101 downward motion in the eye and outside the eyewall. Although this schematic is convenient
102 for describing the basic structure of the inner core in a mature axisymmetric TC, the actual
103 inner core of a TC is not as simple as it has small-scale features and asymmetries, which
104 are dependent on the stage of life cycle, outer core, and environment. In addition, the inner
105 core of the TC can experience a dramatic change such as the Secondary Eyewall Formation
106 (SEF) in which another eyewall develops outside the existing eyewall. The existing eyewall
107 is often replaced with the outer eyewall in a period of one day, which is referred to as the
108 eyewall replacement cycle (ERC). As many researchers have become interested in these
109 topics, new perspectives have been introduced over the last several years, which build upon
110 existing studies. In addition to research purposes, the inner core structure of TCs is also
111 important for operational analyses and forecasting of TC intensity and size. Operations on
112 TC inner core structure in various agencies have recently become more active, however
113 global activities are less documented in publications. Climatological changes in the TC inner
114 core structure are also relevant. Although the development of a high-resolution model that
115 can resolve the TC inner core is expensive for climatological simulation, some recent studies
116 have discussed the structural changes in TC inner cores in future climate scenarios.

117 As numerous studies on TC inner core structure changes have been performed in
118 the last several years, here we propose a review of these activities. The goal is to assist in

119 organizing the available information, motivate further research, and contribute to disaster
120 prevention and mitigation. The paper is organized as follows: In Section 2, a review of recent
121 studies on fine-scale features and the fundamental understanding of the TC inner core,
122 excluding the SEF and ERC, which are described in Section 3. Sections 2 and 3 begin with
123 basic TC dynamics and conventional theories, followed by the recent studies. Recent
124 activities on the operational analyses and forecasting of the TC inner core structure are
125 summarized in Section 4, and the climatological impact is outlined in Section 5. Finally,
126 Section 6 summarizes the conclusions. The review largely focused on literature published
127 since 2018, and built upon existing reviews of the TC inner core structure (Stern et al. 2014a;
128 Vigh et al. 2018). However, some studies published before 2018 were also referenced for
129 better understanding. We do not review the structural changes in the tropical and
130 extratropical transition, which is nicely summarized in Wood et al. (2023). Also, the inner
131 core structure change is strongly related to the TC intensity change. A reader may refer to
132 Chen et al. (2023) for recent studies on TC intensity change.

133

134 **2. TC inner core structure (excluding SEF/ERC)**

135 The basic structure of the TC inner core is reviewed prior to introducing recent
136 progress in TC inner core structure research. A storm undergoes explosive convection,
137 known as a convective burst, during its genesis and development. It has a deep structure
138 that covers the entire troposphere. A mature TC has an eye characterized by a warm core
139 and dry air except near the surface (Hawkins and Rubsam 1968). The eye is surrounded by
140 an eyewall consisting of convective clouds sloping upward and outward, similar to a stadium.
141 The inflow in the boundary layer directs humid air to the eyewall and obtains further moisture
142 from warm sea surface, whereas most of the air flows outward in a shallow layer at the
143 tropopause. The eyewall region exhibited intense diabatic heating and precipitation, and the
144 cloud-free eye was thought to have been formed to compensate for subsidence. The TC

145 tangential velocity is usually maximized at altitudes of 500-1000 m (Powell et al. 2003) and
 146 gradually weakens with increasing height. The tangential wind component that consists of
 147 strong cyclonic circulation is often called the primary circulation, while the radial and vertical
 148 wind components in the radial-pressure section are called the secondary circulation.

149 A mature TC is typically regarded as an axisymmetric vortex in the vicinity of gradient
 150 wind and hydrostatic balanced state (Ooyama 1969). The combination of the gradient wind
 151 balance and hydrostatic balance yields the thermal-wind balance, linking the vertical
 152 decrease in tangential wind to the radial decrease in temperature owing to the warm core.
 153 As the saturated equivalent potential temperature θ_e^* (or moist entropy $s^* = C_p \log \theta_e^*$; C_p is
 154 the specific heat capacity at constant pressure) is nearly conserved above the boundary
 155 layer, the motion of a parcel in the eyewall follows the constant θ_e^* surface. From a dynamical
 156 point of view, the absolute angular momentum (AAM) defined by the following equation is
 157 also an important quantity that is conserved without surface friction or turbulent mixing.

$$M = \frac{1}{2} fr^2 + rv \quad (1)$$

158 where M is the AAM, f is the Coriolis parameter, r is the distance from the TC center, and v
 159 is the tangential velocity. The coincident surface of constant M and constant θ_e^* in the
 160 eyewall region indicates moist slantwise neutrality (Emanuel 1983).

161 The horizontal motion of an air parcel in a rapidly rotating environment is restricted
 162 by inertial stability I^2 defined as follows:

$$I^2 = \left(f + \frac{2v}{r} \right) (f + \zeta) = \frac{1}{2r^3} \frac{\partial M^2}{\partial r} \quad (2)$$

163 where ζ is the vertical component of relative vorticity. The horizontal motion in a rapidly
 164 rotating fluid is constrained by I instead of the Coriolis parameter f as the Rossby number
 165 (v/fr) is much larger than unity. For example, characteristic scales of inertial period and

166 dynamically determined boundary layer depth are $2\pi/I$ and $\sqrt{K/I}$ instead of $2\pi/f$ and
167 $\sqrt{K/f}$, respectively, where K is the diffusive eddy viscosity (Ito et al. 2011; Kepert 2001;
168 Zhang et al. 2011).

169 A mature TC is rather symmetric; however, it also has notable asymmetric features.
170 One of the examples is small scale vorticities called as “mesovorticities” which behave as
171 vortex Rossby waves whose background field is a TC-scale vortex (Macdonald 1968;
172 Montgomery and Kallenbach 1997). Such asymmetric vorticity perturbations can be
173 amplified by barotropic instability and/or forcing. The vertical shear of environmental
174 horizontal wind (simply referred to as VWS) is known to yield an azimuthal wavenumber-1
175 component, in which active convection, heavy rainfall, and local maximum wind speed tend
176 to be observed in the downshear or downshear-left (Corbosiero and Molinari 2003; Kepert
177 2001; Ueno and Kunii 2009). Asymmetries also develop due to the translation speed of a
178 TC. The addition of the TC translation and primary circulation wind vectors results in faster
179 earth-relative winds to the right of the TC motion and enhanced boundary layer radial
180 convergence generally ahead of and to the right of motion (Kepert 2001; Kepert and Wang
181 2001; Shapiro 1983). These asymmetries can project onto the asymmetries created by the
182 VWS (Uhlhorn et al. 2014; Zhang and Uhlhorn 2012).

183

184 *2.1 TC inner-core structure in different stages*

185 While the typical inner core structure of a mature TC is described above, Tao and
186 Zhang (2019) conducted idealized simulations to investigate the structural changes before,
187 during, and after the TC rapid intensification (RI) under various VWS conditions. Prior to the
188 onset of RI, the vortex and convection structure exhibit significant asymmetry. It causes the
189 upper-level warm anomaly induced by diabatic heating to be displaced from the mid-level
190 warm anomaly. Strong updrafts are primarily located on the down-tilt side. Boundary layer

191 inflows positioned underneath and upwind of the updrafts, while outflows are found on the
192 downwind side. On the RI onset, updrafts and downdrafts begin to spread out alternately,
193 encircle the surface center, and overlap. It leads to a temporary weakening of convection.
194 The previously displaced warm cores start to merge, and the radial flows in the boundary
195 layer become more axisymmetric, consistent with findings by Miyamoto and Takemi (2013).
196 During RI, although some asymmetry remains, the structure of an eye became more upright,
197 and the eyewall had strong updrafts with downdrafts outside the eyewall and subsidence in
198 the eye. These processes are illustrated in Fig. 1. Alvey III and Hazelton (2022) investigated
199 the misaligned vortices located in the middle and lower troposphere can align due to deep
200 convection (not necessarily symmetric). Using ensemble simulations, Alvey et al. (2020)
201 showed that the RI members are characterized by persistent and deep convection in the
202 downshear. As a result, the vortex became aligned through the stretching term and
203 precession. The stratiform precipitation and anvil clouds were seen in the upshear quadrants.
204 Condensate transported from the downshear quadrants makes the middle to upper
205 tropospheric air more humid by evaporation and sublimation. This contributes to the
206 symmetrization of precipitation region, which is a necessary precursor for RI. Chen et al.
207 (2018) and Chen et al. (2021) displayed that higher sea surface temperatures (SSTs) can
208 facilitate more symmetric inner core precipitation which promotes faster vertical vortex
209 alignment and faster contraction of the radius of maximum wind (RMW) preceding RI.

210 Even in the axisymmetric model, the inner core structure differs significantly
211 between the developing and mature stages. Peng et al. (2019) displayed how the constant
212 AAM and θ_e^* surfaces evolved to become parallel from an almost orthogonal state in the
213 early development stage (Fig. 2). In the early development stage (Phase I), sporadic and
214 deep convection from the high s^* air near the sea surface repeatedly redistribute large M
215 and s^* values vertically from the lower troposphere to the tropopause; the locally large value

216 of M is associated with convective rings. In the later development stage (Phase II), the
217 congruent M and s surfaces coevolve toward the center and become aligned. Fei et al.
218 (2021) investigated the role of vertical advection on the boundary layer structure,
219 intensification rate, and intensity. The removal of vertical advection of radial wind
220 significantly reduced both the height and strength of supergradient wind core. Meanwhile,
221 the removal of vertical advection of agradient wind reduced the height of supergradient wind
222 core but slightly increased the strength.

223 An observational study on Hurricane Michael (2018) confirmed that M surfaces are
224 redistributed in the inner core region and that RI is accompanied by the $\partial z / \partial r$ increase
225 along the constant M surface in the inner core (DesRosiers et al. 2022). It supports the
226 previous study showing the smaller RMW with the larger $\partial z / \partial r$ (Stern et al. 2014b).
227 Meanwhile, Ito and Yamamoto (2022) argued for the slope of a constant M surface using
228 commercial aircraft data for TCs, most of which were in the decaying stage in the western
229 North Pacific. Their analysis showed that strong TCs tended to have a small $\partial z / \partial r$ in the
230 upper troposphere. Considering that the slope corresponds to the inertial stability divided by
231 the baroclinicity, they believed that changes can be attributed to the strong potential
232 temperature gradient of strong TCs.

233 Other environmental factors, such as VWS, synoptic-scale atmospheric circulation,
234 and upper-level troughs, also influence the inner core structure. Using a database of
235 composited airborne Doppler radars, Wadler et al. (2018) found that intensifying TCs have
236 updrafts at higher altitudes and stronger magnitudes in the upshear quadrants than steady-
237 state TCs. Strong and deep updrafts located inside the RMW efficiently spun the vortex
238 (Rogers et al. 2013; Zhang and Rogers 2019).

239 Li et al. (2019) revisited the dynamics of eyewall contraction. Using the azimuthal-
240 mean tangential wind budget and the diagnostic equation proposed by Stern et al. (2015),
241 in which the change of RMW is decomposed into the radial gradient of local time tendency

242 of the azimuthal mean tangential wind and the curvature of the radial profile of azimuthal-
243 mean tangential wind, they showed radial advection dominated the RMW contraction in the
244 lower boundary layer and the vertical advection causes the RMW contraction in the upper
245 boundary layer and lower troposphere. Also, the eyewall contraction is dependent on the
246 horizontal mixing near the eyewall. Wu and Ruan (2021) reported that the rapid RMW
247 contraction tends to precede the RI in observations. The possible mechanism of this
248 phenomena was given by Li et al. (2021b). First, they affirmed that the AAM is not conserved
249 following the RMW. During the rapid contraction phase, the RMW decreases inward across
250 the AAM surfaces, explaining the RMW decreases with relatively small intensification. On
251 the other hand, when the RMW contraction slows down and the intensification becomes
252 rapid, the AAM surfaces move inward passing through the RMW.

253

254 2.2 Nontraditional features in the inner core

255 When Hurricane Patricia (2015) reached its lifetime maximum intensity, an atypical
256 inner core structure was observed, with a maximum tangential wind in the mid-troposphere,
257 in addition to the conventional maximum at the lower level (Stern et al. 2020). This feature,
258 which is ultimately due to the inertial oscillation driven by surface friction, was found to be
259 prominent in intense and/or small storms. The numerical simulation displayed that the multi-
260 maxima disappeared when the vertical mixing was strong. Yamada et al. (2021) described
261 double warm cores in the eye region of TC Lan (2017) based on dropsonde observations
262 during the Tropical Cyclones–Pacific Asian Research Campaign for the Improvement of
263 Intensity Estimations/Forecasts (T-PARCII) aircraft campaign (Fig. 3). Their saturation point
264 analyses revealed that the air parcels in the upper warm core originated from the eyewall
265 rather than the lower stratosphere. Furthermore, the dropsondes indicated the vigorous
266 turbulent mixing between eye and eyewall. Based on these findings, they proposed a
267 conceptual model in which the parcels in the eye boundary layer were mixed through eddies

268 and moved upward on the inner side of the eyewall updrafts and consisted of the upper-
269 level warm core. Double warm cores have been reproduced by simulations (Stern and
270 Zhang 2013) and Hurricane Patricia also appeared to have a double warm core structure
271 based on high density Global Hawk dropsonde data (Rogers et al. 2017).

272 Inside the eye region, sporadic short-lived deep convective clouds of up to 9 km,
273 which were considerably taller than stratocumulus hub clouds of up to 3 km (Simpson
274 1952), were observed in the decay stage of TC Trami (2018) during T-PARCII. Both
275 observations and numerical simulations display that these clouds formed after the
276 weakening and outward shift of the eyewall due to significant sea-surface cooling because
277 of the slow translation speed of TC Trami (2018) (Hirano et al. 2022). Their analysis
278 showed that the genesis of deep eye clouds may indicate the decay of the warm core in
279 the middle-lower troposphere.

280 The striating clouds (“striations” or “finger-like features”) at $r \sim 30$ km were clearly
281 observed by the Himawari-8 satellite, which had remarkably high angular velocities at the
282 inward flank of the eyewall (Fig. 4). Their radial orientation is gradually tilted in a time scale
283 of 10 minutes according to the differences in the tangential speed depending on the altitude
284 (Tsukada and Horinouchi 2020).

285

286 2.3 Wavenumber-1 component

287 Geostationary satellite Himawari-8 has a special mode to take images every 30 seconds.
288 Horinouchi et al. (2023) utilized those rapid-scan observations for analyzing the flows in the
289 eye region of TC Haishen (2020) during the intensification. The low-level local circulation
290 center was located several kilometers away from the storm-scale center. Their analysis
291 showed that the local circulation behaves wavenumber-1 component which developed
292 algebraically (Nolan et al. 2001). An important role of the local circulation is to homogenize
293 the rotational angular velocity within the eye and contributes to the increase of azimuthal-

294 mean tangential wind. Dai et al. (2021) found that the wavenumber-1 convection asymmetry
295 in the outer eyewall of Typhoon Lekima (2019) was not integrally related to the ambient
296 VWS or frictional effect associated with TC motion as in (Kepert 2001). They suggested that
297 this asymmetry results from the phase locking between the radially-outward propagating
298 wavenumber-2 VRWs, which originate from the inner eyewall, and the wavenumber-1
299 asymmetry that cyclonically propagates near the radially-inward side of the outer eyewall.

300

301 *2.4 Boundary layer modulation on TC inner core structure*

302 Ahern et al. (2019) analyzed dropsonde composites of the TC boundary layer
303 structure at different stages of the TC lifecycle. Intensifying TCs are characterized by deeper
304 jets above the logarithmic layer in the eyewall region, higher (weaker) inertial stability within
305 (the outer area of) the eyewall, and more inward-penetrating boundary layer radial inflow
306 than weakening TCs. Focusing on the decay phase of a full-physics simulation of Hurricane
307 Irma (2017), Ahern et al. (2021) analyzed how the asymmetric agradient forcing within the
308 boundary layer from the VWS and the storm motion force local accelerations of the primary
309 vortex, suggesting that the interplay between the directionality of the VWS and storm motion
310 vectors produces structural asymmetries. Zhang et al. (2023) analyzed a large number of
311 Doppler profile data in multiple storms and derived composites of boundary layer structure
312 of TC with different intensity and intensity change rates. First, they showed that the depth of
313 an inflow layer is deeper in weaker storms than in strong storms. They also found that
314 intensifying TCs tend to have the rapid radial decay of the tangential wind outside of RMW
315 (called “narrow”). The narrow vortex is also characterized with larger inflow and strong
316 updrafts being located near RMW, while non-intensifying TCs have large inflow in the outer-
317 core region. The total advection of mean AAM near the RMW is much larger in narrow vortex
318 than in broad vortex suggesting larger intensification rate given the same intensity of these
319 two types of vortices.

320 While VWS is typically not favorable for TC intensification, Hurricane Michael (2018)
321 intensified under moderate VWS. Using dropsondes and idealized simulations, Wadler et al.
322 (2021a) and Wadler et al. (2021b) found that enhanced low-level outflow upshear originating
323 from the inner-core insulated the boundary layer inflow from low-entropy convective
324 downdrafts, aiding in boundary layer flux recovery. Rios-Berrios et al. (2018) also used
325 idealized simulations to explore intensification under moderate VWS. Their analysis
326 revealed that a reduction in tilt and increased symmetry precede intensification, as these
327 processes are linked to a significant increase of near-surface vertical mass fluxes and
328 equivalent potential temperature. A vorticity equation indicated that the increase of near-
329 surface vertical mass fluxes facilitate intensification by stretching near the surface and tilting
330 of horizontal vorticity in the free troposphere. Notably, the reduction in tilt occurs due to
331 vortex merger that forms a single closed circulation. This vortex merger occurs after the
332 vortex in the middle troposphere moves to the upshear left. It means that the flow supports
333 near-surface vortex stretching, deep updrafts, and the reduction of low-entropy fluxes in the
334 upshear left quadrant. Zhang et al. (2013) showed that anomalously large enthalpy fluxes
335 located to the left of the shear vector, owing to the presence of mesoscale ocean eddies,
336 also aided in faster boundary layer recovery. It is supported by Chen et al. (2018) and Chen
337 et al. (2021) indicating that higher SST has been found to aid symmetrization through faster
338 boundary layer recovery. In addition, strong upper ocean salinity vertical gradients can
339 increase the upper ocean density stratification and reduce SST cooling (Hlywiak and Nolan
340 2019; Rudzin et al. 2018). Le Hénaff et al. (2021) and John et al. (2023) showed that
341 anomalously fresh sea surface salinity (SSS) due to enhanced outflow from the Mississippi
342 River reduced SST cooling under Hurricanes Michael and Sally, respectively, which
343 facilitated axisymmetrization. Strength of the vertical turbulent mixing also regulates the
344 boundary layer inflow and recovery process through multiscale interaction (Zhang and
345 Rogers 2019; Zhang et al. 2017b).

346 Chen and Chavas (2020) and Hlywiak and Nolan (2021) investigated the time-
347 dependent response of a TC structure to surface roughening and drying, analogous to
348 landfall scenarios, based on idealized simulations. Chen and Chavas (2020) studied the
349 axisymmetric response to instantaneous changes in the surface, whereas Hlywiak and
350 Nolan (2021) simulated the three-dimensional response of translating systems. Both studies
351 found similar results: surface drying produced monotonic weakening and horizontal
352 expansion of the vortex as reduced moisture fluxes into the BL eventually stabilized the
353 eyewall and weakened secondary circulation. The response to drying was initially weak,
354 however became more apparent over time after significant weakening occurred. Higher
355 surface friction produces a more immediate effect compared to drying through the
356 deceleration of the primary circulation and weakening of inertial stability. This supports the
357 transient acceleration of the secondary circulation, enhancing the radial inflow and angular
358 momentum import within the BL and accelerating the outflow above the boundary layer
359 before weakening and inner core expansion commenced. Subsequently, Chen and Chavas
360 (2020) found that the total response of the maximum winds is approximately a product of
361 the individual responses to surface roughening and drying, indicating the applicability of the
362 maximum potential intensity theory (Emanuel 1986) to the landfall decay process. Hlywiak
363 and Nolan (2022) followed up Hlywiak and Nolan (2021) with an investigation of the
364 asymmetric response of the inner-core boundary layer during landfall and found that the
365 accelerated secondary circulation response to increased friction begins over land and within
366 the offshore flow preceding landfall of the eye. Friction over land forces a deep subgradient
367 inflow layer, subsequently enhancing the near-eyewall convergence within the offshore flow
368 and accelerating the downstream supergradient boundary layer jet (Fig. 5). These
369 asymmetries in the inner core boundary layer at landfall are due to the sharp gradient of
370 surface frictional stress, a mechanism which is similar yet larger in magnitude to the
371 asymmetries that develop for translating TC over the open ocean (Kepert and Wang 2001;

372 Shapiro 1983). Aircraft observations of landfalling Ida (2021) documented boundary layer
373 inflow and wind asymmetry from ocean to land (Rogers and Zhang 2023) generally
374 supporting Hlywiak and Nolan's findings in numerical simulations.

375 TC boundary layers contain a lot of small-scale coherent structures, such as roll
376 vortices and tornado-like vortices (e.g. Ito et al. 2017; Zhang et al. 2008). Recent LES
377 simulations have demonstrated that a boundary layer becomes shallower with stronger VWS
378 near the surface by employing the fine-mesh. For example, Ito et al. (2017) identified three
379 types of near-surface coherent structures. A Type-A roll nearly parallel to the tangential wind
380 prevailed outside of the RMW, and Type-B rolls almost orthogonal to the type-A rolls were
381 found near the RMW. They were caused by an inflection-instability. Type-C rolls nearly
382 parallel to the tangential wind were found inside the RMW, which can be explained by a
383 parallel instability. When the grid spacing is 100 m or less, roll vortices significantly contribute
384 to the transfer of heat and momentum (Liu et al. 2021). Similarly, LES simulations of a
385 landfalling TC have replicated the presence of a thin boundary layer and highlighted the
386 critical role of roll vortices (Li et al. 2021a).

387

388 *2.5 Dry and moist TCs*

389 Several studies have attempted to establish a fundamental understanding of the
390 contribution of moisture to the inner core structure of TCs, including the radius of maximum
391 wind (RMW), width of the eyewall, and depth of the inflow layer. Generally, TC-like storms
392 (hereafter, dry TCs) are known to potentially exist by relying only on sensible heat fluxes
393 without water substances (Mrowiec et al. 2011), although real TCs are tied to water in Earth's
394 present climate. By comparing typical TCs to dry TCs, Cronin and Chavas (2019)
395 investigated the role of moisture on TC structure. The RMW will become substantially larger
396 when a TC is maintained under dry conditions, and the inflow layers of the dry TCs are much
397 deeper than those of typical TCs. Wang and Lin (2020) further investigated the differences

398 between moist and dry TCs using idealized simulations and found that the sub-saturation of
399 a moist atmosphere, not the moisture itself, causes distinctions between moist and dry TCs.
400 The inner core structures of the moist reversible and dry TC were almost the same, whereas
401 a typical moist TC with falling hydrometeors was significantly different from a dry or moist
402 reversible TC (Fig. 6). It suggests that falling hydrometeors place a strong constraint on the
403 inner core structure of TCs through redistributing entropy. However, the mechanism by
404 which the irreversible entropy production modulates the inner core structure of TCs remains
405 unknown.

406

407 *2.6 Impact of TC outer core and environment on storm structure*

408 The outer core structure was also found to be an important factor influencing the
409 inner core structure. Using an axisymmetric nonhydrostatic model, Tao et al. (2020) revisited
410 Lilly's model, which relates the tangential wind structure of an axisymmetric steady-state TC
411 to its outer size and environmental conditions by simplifying the assumptions. They showed
412 that variations in the sea surface, boundary layer and tropopause temperatures, and the
413 AAM at some outer radius affect the radial profile of the tangential velocity at the top of the
414 boundary layer. On the other hand, these parameter variations do not substantially affect
415 normalized tangential velocity, $V(r/r_m)/V_m$, where V_m is the maximum tangential wind at
416 RMW. Chavas et al. (2015) demonstrated that the RMW could be reasonably estimated from
417 the outer size by applying a TC structure model. They found that the eyewall expansion of
418 Hurricane Helene (2006) during its RI was due to environmental-scale low-level
419 convergence and upper-level divergence of the angular momentum flux by accelerating the
420 tangential velocity outside the eyewall.

421 Martinez et al. (2020) found that initially large TCs remained large during
422 development by importing a large angular momentum into the inner core and spinning up

423 the wind field (Fig. 7). Martinez et al. (2020) also displayed that environmental factors such
424 as the amount of moisture can modulate the outer core structure, which in turn has a
425 significant effect on the inner core structure. The moist environment promotes outer core
426 convection, spins up the outer core wind field, and leads to the contraction of large angular
427 momentum surfaces into the inner core of the TC, which expands the horizontal field of the
428 tangential velocity (Fig. 7). Qin and Wu (2021) also confirmed this environmental impact on
429 the outer core structure and the consequent impact on the inner core structure. Shen et al.
430 (2021) revealed that limiting the enthalpy fluxes within the outer core reduces the total inward
431 transport of angular momentum, leading to a smaller inner core size, whereas suppressing
432 the fluxes within the inner core has little influence on the overall storm size. They generally
433 support the findings of Xu and Wang (2010b) and Xu and Wang (2010a), which investigated
434 the relationship between an inner core and an outer core and impact of environmental
435 moisture. These results highlight the importance of the spatial heterogeneity of surface
436 fluxes in the inner core boundary layer structure.

437

438 **3 SEF and ERC**

439 TCs often possess a secondary eyewall outside the original eyewall (Willoughby et
440 al. 1982) as illustrated in Fig. 8, and SEF events are more likely to occur as the TCs intensify
441 (Kossin and Sitkowski 2009; Kuo et al. 2009). It typically begins with a spiral rain band
442 morphing into a new eyewall outside the inner eyewall. Kuo et al. (2009) reported that
443 approximately 57% and 72% of Category 4 and Category 5 TCs, respectively, possessed
444 concentric eyewalls at some point during their lifetime in the western North Pacific. After the
445 secondary (or outer) eyewall is generated, the inner core region exhibits multiple maxima of
446 tangential winds and precipitation (Sitkowski et al. 2011). The low-echo reflectivity region
447 between the concentric eyewalls is called a moat. Following the SEF, TCs often undergo an
448 ERC, characterized by the weakening of the original inner eyewall and its replacement by

449 the outer eyewall, which becomes the new primary eyewall. If the environmental conditions
450 remain favorable for TC development, TCs can then reintensity. Such a process is typically
451 over a timeframe ranging from a couple of hours to a day or two (Kossin and DeMaria 2016;
452 Sitkowski et al. 2011). There are also cases in which SEF TCs do not undergo an ERC;
453 instead, the outer eyewall decays and the inner eyewall retains (Yang et al. 2013). Although
454 the inner eyewall typically decays and the outer eyewall contracts with intensification on a
455 timescale of less than a day or two, multiple eyewalls can sometimes coexist for a longer
456 period, or the outer eyewall decays earlier (Kossin and DeMaria 2016; Yang et al. 2013).
457 Many physical mechanisms have been proposed for the SEF and ERC (as summarized in
458 Stern et al. 2014a; Vigh et al. 2018), yet they are not fully understood.

459 Simplified models are known to reproduce the structure similar to a secondary
460 eyewall. Nong and Emanuel (2003) and Emanuel et al. (2004) reported that an SEF-like
461 structure is generated even with an axisymmetric model. Whereas SEF events are
462 reproduced much less in an axisymmetric model, it may aid in the exploration of certain
463 aspects of plausible mechanisms. Using a two-dimensional non-divergent barotropic model,
464 a strong inner vortex representing the TC core can envelop neighboring weak vorticity
465 patches, forming an SEF-like vorticity ring (Kuo et al. 2004). However, Moon et al. (2010)
466 argued that the allocation of convection-induced small vorticity dipoles, which have
467 comparable strength to the inner-core vorticity and potentially better represent moist
468 convective vorticity in the TC's outer core, did not result in coherent concentric vorticity rings
469 in barotropic two-dimensional flow. In contrast, Moon and Nolan (2010) successfully
470 simulated a secondary horizontal wind maximum by prescribing the same vorticity dipoles
471 using a three-dimensional model that calculates more sophisticated physical processes.
472 These findings underscore the importance of considering the complexity of three-
473 dimensional processes in capturing SEF TCs.

474 The SEF was associated with the upward motion outside the primary eyewall. In

475 addition to subsidence induced by the updraft in the inner eyewall, Rozoff et al. (2006)
476 proposed that rapid filamentation dominated by strain plays a crucial role in moat formation.
477 This rapid filamentation distorts vorticity generation within a 30-minute timescale, impeding
478 deep convection. Observations indicate that mature hurricanes typically exhibit a radial area
479 with a slightly negative radial gradient of azimuthal-mean vorticity beyond the RMW at lower
480 levels, known as a beta-skirt region (Mallen et al. 2005). Within this region, under conditions
481 of rapid filamentation, sufficient convective available potential energy (CAPE), and low
482 convective inhibition, sporadic convection can persistently occur. Building upon the
483 framework of two-dimensional theory (Vallis and Maltrud, 1993), Terwey and Montgomery
484 (2008) argued that this convective energy likely upscales towards the tangential direction
485 over the beta-skirt region, leading to the accumulation of a ring of low-level tangential jets
486 outside the primary eyewall.

487 Energy in the SEF region can accumulate through the stagnation of vortex Rossby
488 waves (Montgomery and Kallenbach 1997), which is related to the slowdown of the outward
489 group velocity with increasing radial wavenumbers in the outer region. Some studies support
490 the contribution of the accumulation of eddy kinetic energy near the stagnation radius to the
491 SEF (e.g. Menelaou et al. 2012), while some studies do not (e.g. Qiu et al. 2010). In addition,
492 humid outer environments can be one of the important factors for the SEF (Ge 2015). This
493 is because the unstable condition is favorable for TC to establish a secondary eyewall in the
494 outer region.

495 In the recent decades, numerical models have demonstrated the increasing ability in
496 reproducing the SEF occurrence; however, there remains significant uncertainty in
497 simulated SEF regarding the onset time and duration of ERCs, as well as in predicting
498 changes in TC intensity and structure during ERCs. Three-dimensional high-resolution
499 models with a horizontal grid spacing of less than a few kilometers are needed for
500 reproducing the SEF and ERC (e.g., Houze et al. 2007). Analyzing the simulated Typhoon

501 Sinlaku (2008), Wu et al. (2012) and Huang et al. (2012) identified a sequence of structural
502 changes within and just above the boundary layer, and proposed an unbalanced dynamical
503 pathway to SEF, which was further elaborated by Huang et al. (2018). Some elements of
504 this unbalanced dynamical pathway were also revealed by numerical modeling data (Abarca
505 and Montgomery 2013) and data collected from flight observations (e.g. Abarca et al. 2016;
506 Didlake and Houze 2011). Taking a different perspective, Kepert (2013) proposed a formula
507 showing the analytical solution for the frictionally-induced boundary layer updraft, and
508 emphasized the coupling and positive feedbacks between the friction-induced updraft and
509 the aloft convection. The hypothesis of Kepert (2013) was supported by Kepert and Nolan
510 (2014), Zhang et al. (2017a), Kepert (2017), and Kepert (2018).

511 As the tangential wind field expands preceding SEF, a corresponding augmentation
512 in surface heat fluxes beyond the eyewall is anticipated. Cheng and Wu (2018) conducted
513 numerical experiments where they applied various caps on surface winds, used for surface
514 flux calculations, across different radial intervals, illustrating the importance of the WISHE
515 mechanism in SEF and ERC. In these numerical experiments, SEF failed to occur when
516 surface heat fluxes were markedly suppressed both inside and outside the SEF region.
517 When surface heat fluxes were moderately suppressed within the same area, SEF initiation
518 was delayed, and the intensity of both eyewalls diminished. Notably, the suppression of
519 surface heat fluxes in the inner-core region had negligible impact on the outer eyewall
520 evolution. Observations and numerical experiments have shown that SEF tends to occur
521 under weak-to-moderate VWS conditions (Didlake et al. 2017; Zhang et al. 2017a). It implies
522 the contribution of asymmetric features, such as a rainband. Other factors are known to be
523 related to SEF. For example, no-ice sensitivity experiments displayed that the SEF was
524 delayed without an ice phase (Terwey and Montgomery 2008).

525 The aforementioned mechanisms and speculations are not necessarily exclusive and
526 may reflect different aspects of the same phenomenon. Alternatively, some characteristics

527 may only be valid in certain special cases. The SEF and ERC are topics that have been
528 continuously investigated. Very recent studies have covered the above-mentioned
529 mechanisms in more detail, display evidence in real case observations, propose a new
530 theory, and investigate special cases. Additionally, the decay of the inner eyewall was
531 investigated. These advances are summarized in the following subsections.

532

533 *3.1 Classification of SEF and ERC*

534 Cheung et al. (2024) investigated a dataset of 87 secondary eyewall progressions
535 primarily based on 89-92-GHz passive microwave imagery. In their works, the secondary
536 eyewall is defined as the azimuthal coverage of at least 50% surrounding the inner
537 eyewall. They classified the first stage into two types: a spiral band stemming from the
538 inner eyewall and an outer rainband. The former typically progressed into half-concentric
539 coverage, and the latter progressed into the full-concentric coverage. The exit stage
540 consists of two distinct categories: Replacement or no replacement. The “no replacement”
541 can be further categorized into three types: the secondary eyewall fading on microwave
542 imagery, the merging of the secondary eyewall with the inner eyewall and becoming a
543 spiral rainband after the loss of concentricity. One common pathway is that an outer
544 rainband develops into a fully concentric eyewall followed by an ERC, while another typical
545 pathway involves an inner rainband forming a half-concentric eyewall without any
546 replacement. ERC events favor more intense storms, weaker VWS and greater relative
547 humidity.

548

549 *3.2 Stationary band complex (SBC) and SEF*

550 Based on airborne Doppler radar analysis, Didlake et al. (2018) argued that the SEF
551 was initiated from an organized rainband complex comprising convective precipitation in the
552 downshear-right and stratiform precipitation in the downshear-left to upshear-left (Fig. 9).

553 This rainband complex remained quasi-stationary relative to the storm center, like the SBC
554 discussed in Willoughby et al. (1984). In the downwind sections (left of the VWS) of the
555 rainband complex, falling ice particles formed a wide and uniform precipitation band with
556 predominantly stratiform characteristics. Stratiform precipitation in the SBC induced
557 mesoscale descending inflow (MDI). This process aided in the expansion of the vortex
558 circulation and transported mid-level low-equivalent potential temperature (θ_e) into the
559 boundary layer. The updraft next to the MDI was seemingly caused by convergence and
560 upward acceleration. It was induced by the negative buoyancy of MDI, entering the high- θ_e
561 region in the boundary layer. This updraft and the MDI in the downshear-left quadrant
562 increased the tangential velocity, and it yielded the axisymmetric structure of the secondary
563 eyewall. This can explain why SEF events typically occur under weak-to-moderate VWS
564 and why SEF is sensitive to the ice phases.

565 Vaughan et al. (2020) compiled a five-year climatology of the SBC based on passive
566 microwave satellite data. They demonstrated that approximately 80% of 84 SEF events were
567 preceded by an SBC in the 6 h time window prior to SEF. The geometry of the SBC has
568 higher azimuthal extent and lower crossing angle from 12–24 h before the SEF (Fig. 10).
569 This suggests that SBC-associated dynamic processes are important for SEF. However,
570 SBCs with substantial radial extent and high circularity without subsequent SEF were also
571 detected in the dataset (Fig. 10c).

572 Yu and Didlake (2019) used idealized simulations with a full-physics three-
573 dimensional model to investigate the response to a prescribed heating that mimicked the
574 contribution of stratiform around spiral TC rainbands, similar to Moon and Nolan (2010). The
575 heating structure represented a quasi-stationary rainband complex. The vortex response
576 included the MDI and a low-level forced radial updraft inside the rainband heating.
577 Stratiform-induced cold pool interacted with tangential flow of upwind warm air, and it caused
578 the updraft through buoyancy. It clearly shows that the diabatic forcing of rainband is critical

579 for triggering and sustaining the forced low-level updraft, which might lead to a SEF.

580 Yu et al. (2021a) and Yu et al. (2022) further investigated the role of the stratiform
581 rainband processes using a full-physics simulation of Hurricane Matthew (2016). Yu et al.
582 (2021a) focused on the tangential velocity in the pre-SEF stage. Under a moderate VWS,
583 the storm developed a quasi-SBC that aligned with the shear vector. Prior to the SEF, the
584 storm experienced a broadening of the tangential velocity, which was largely by horizontal
585 advection of momentum according to MDI in the downshear-left stratiform region. The MDI
586 was connected to the boundary layer in the upshear-left quadrant, and convections were
587 induced along its inner edge. It helped develop the maximum of low-level tangential velocity
588 within the incipient secondary eyewall. Yu et al. (2022) focused on the emergence,
589 maintenance, and impact of these persistent updrafts which developed within the left-of-
590 shear quasi-SBC. Updraft initiation was based on the process described by Yu and Didlake
591 (2019). In this process, buoyancy advection resulted from an MDI-induced cold pool
592 interacting with the high- θ_e air in the inner core. A budget analysis of θ_e demonstrated that
593 the updrafts in the left of the VWS were maintained by enhanced moist instability due to
594 differential horizontal advection in the boundary layer. A potential vorticity (PV) budget
595 showed that these updrafts generated PV anomalies that propagated cyclonically downwind.
596 The propagation of PV anomalies downstream increased the azimuthal mean PV at the
597 radius of the SEF, corresponding to the axisymmetrization of the secondary eyewall itself.
598 These studies provide a dynamical explanation, as illustrated in Fig. 9, of how stratiform
599 rainband processes can explain the SEF in a sheared, mature TC. Recent studies, such as
600 those by Wang and Tan (2020), Zhu et al. (2022), and Kasami and Satoh (2024)
601 demonstrated a similar MDI pathway to SEF in their modelling studies. Wang et al. (2019a)
602 conducted an idealized simulation under the axisymmetric environment on an f -plane, yet
603 the rainband contributed the simulated SEF in a similar manner.

604 Recent observational evidence of this stratiform pathway to the SEF has also been

605 presented (Fischer et al. 2020; Kanada and Nishii 2023; Razin and Bell 2021). Fischer et al.
606 (2020) examined the RI and associated ERCs of hurricane Irma (2017) using flight-level and
607 airborne radar observations, microwave satellite observations, and model-based
608 environmental analyses. Irma's RI event included two short ERC episodes (less than 12 h).
609 During the first SEF event, the upward motion and tangential velocity had a secondary peak
610 at the forefront of the MDI. Kanada and Nishii (2023) focused on ground radar-based
611 observations of the outer eyewall formation and inner eyewall weakening of TC Hinnamnor
612 when it approached to Okinawa's main island in 2022. They revealed that an SBC developed
613 in the down-to-left quadrant of a moderate VWS that turned into an outer eyewall. As the
614 VWS weakened, the SBC became more axisymmetric, and a weakly convective "moat"
615 region appeared between the outer and inner eyewalls.

616

617 *3.3 boundary layer processes*

618 The importance of boundary layer processes on SEF has been reconfirmed by
619 many studies. Kepert (2018) developed a boundary layer model for TC rainbands. The
620 boundary layer flow yields a strong low-level convergence and consists of an updraft along
621 the rainband and some region located downwind. The upper boundary layer has a marked
622 wind peak along the band that was approximately 20% stronger than the balanced flow.
623 Therefore, the secondary peak of wind at the SEF region can be explained by the
624 boundary layer dynamics near a rainband. Yu et al. (2021b) investigated the SEF using
625 the Kepert (2018) nonlinear boundary layer model and storm composites of tangential
626 wind observations from Wunsch and Didlake (2018). For the pre-SEF composite, the
627 boundary layer model response exhibited clear secondary maxima in the updrafts,
628 tangential wind, and radial inflow in the left-of-shear quadrant, which contributed
629 substantially to the azimuthal mean. This finding suggests that, leading to SEF, the earliest
630 signal of coupling with the boundary layer and free troposphere is likely to occur in these

631 left-of-shear sectors. As this is also the same region where the MDI triggered new updrafts
632 in previous studies on individual storms, these studies support the idea that the rainband
633 complex process detailed in the previously discussed studies is likely a prevalent
634 mechanism for SEF in sheared storms.

635 Huang et al. (2012) and Wu et al. (2012) identified dynamical processes consistent
636 with unbalanced boundary layer features such as supergradient wind, and their
637 contributions to SEF were revealed by the momentum budget of tangential and radial
638 winds (Huang et al. 2018). These findings continue to be demonstrated in many studies
639 (Ahern et al. 2022; Chen 2018; Guan and Ge 2018; Persing and Montgomery 2022; Razin
640 and Bell 2021; Wang et al. 2019a; Wang et al. 2018; Wang and Tan 2020, 2022). Most of
641 these studies emphasized the cooperation of this boundary layer pathway with other
642 mechanisms occurring in an aloft-free atmosphere, such as the asymmetric dynamics
643 associated with outer rainbands.

644 Miyamoto et al. (2018) proposed that the SEF can be attributed to the instability of
645 the flow in a free atmosphere coupled with Ekman pumping. Vertical wind perturbation is
646 damped in the classic Ekman theory, which considers slow flow in a horizontally uniform
647 structure. Meanwhile, a positive feedback mechanism for vertical wind perturbation works
648 in curving fast flows. This instability tends to be satisfied when an angular velocity is high,
649 absolute vertical vorticity is low, radial gradient of angular velocity is small and negative,
650 and gradient of vertical vorticity is strongly negative. In their setup, the unstable condition
651 is satisfied only between two and seven times the RMW. They verified that the necessary
652 condition for this instability was satisfied in the SEF simulated using a full-physics three-
653 dimensional model.

654

655 *3.4 SEF/ERC in a sheared environment*

656 Liu et al. (2022) and Wang and Tan (2022) reproduced the SEF/ERC in simulations

657 under moderate-strong VWS and weak to moderate VWS to a vortex with maximum wind at
658 approximately 45 m s^{-1} and 70 m s^{-1} (or stronger), respectively. They found that increasing
659 the VWS caused earlier SEF through the formation of outer rainbands. Moreover, Wang and
660 Tan (2022) showed that a vortex with a larger outer core size or a stronger outer wind field
661 could undergo ERC in a non-sheared or strongly sheared environment. In a simulation of
662 Hurricane Earl (2010), Ahern et al. (2022) argued that Earl's azimuthal asymmetry in low
663 level wind and thermal fields affects the azimuthal structure of the broadening of low-level
664 swirling winds; therefore, the secondary maximum winds first appear in the downshear-left
665 quadrant. Low-level azimuthal asymmetry was related to moderate-to-strong VWS and
666 asymmetric friction due to TC motion. By analyzing satellite data, Yang et al. (2021)
667 demonstrated that during ERCs, the mean VWS in the environment for long-lived (longer
668 than 20 h) and short-lived ERC events were weak and moderate, respectively.

669

670 *3.5 Development of warm-core ring and moat prior to SEF*

671 In an idealized numerical simulation, Wang et al. (2019b) found that an off-center
672 warm ring in the upper level outside the eye emerged prior to the SEF and during the
673 broadening of the tangential wind field, and then rapidly strengthened after the SEF. They
674 suggested that this off-center warm ring is a plausible indication of the subsequent
675 occurrence of SEF. Qin et al. (2021) performed a semi-idealized simulation of Typhoon
676 Matsa (2005) and found that the moat formed with a well-developed anvil cloud extending
677 outward from the eyewall. Beneath the anvil cloud, the heating-induced inflow appeared to
678 enhance sublimation cooling by drying local conditions. Subsidence is further enhanced by
679 the downward motion that compensates the strong updraft in the eyewall. Furthermore, the
680 strong filamentation effect and reduced ambient VWS facilitated the axi-symmetrization of
681 the moat and SEF. The authors suggested that a strong eyewall with a well-developed anvil
682 in strength and space may favor SEF when the VWS environment transitions from moderate

683 to weak magnitude. This feature is supported by an observational study (Kanada and Nishii
684 2023).

685

686 *3.6 Role of VRWs*

687 Several studies have investigated the role of vortex Rossby wave activity and
688 examined its role in the SEF or intensity/structure change during an ERC process. Fischer
689 et al. (2020) indicated that the theoretical values of VRW's stagnation radius, provided by
690 Montgomery and Kallenbach (1997), are close to the SEF regions of Hurricane Irma (2017).
691 Guimond et al. (2020) presented evidence of VRW activity in the outer eyewall of Hurricane
692 Matthew (2016) under moderate to strong VWS. After ERC completion, Hurricane Matthew
693 did not reintensify as it contracted, in addition, it was not substantially weakened under the
694 high shear environment. Cha et al. (2021) investigated the ERC process of Hurricane
695 Matthew using ground-based and airborne radars. During the ERC, the VRW damping
696 mechanism proposed by Reasor et al. (2004) could assist the vortex resiliency and
697 resistance to increasing VWS. The AAM budget analysis by Guimond et al. (2020) also
698 indicated that VRW dynamics appear to build up a second SEF event. Flight-level
699 observations of Hurricane Matthew (2016) displayed secondary peaks of tangential velocity
700 where VRW were active, implying their connection to the SEF process.

701

702 *3.7 Barotropic and baroclinic instability during the ERC*

703 Observations have shown that the location of inner eyewall often wobbles with a
704 cycloidal track when multiple eyewalls form. Menelaou et al. (2018) demonstrated that the
705 wobbling comes from the wavenumber-1 instability that grows exponentially in a three-
706 dimensional framework (Fig. 11). This instability can be interpreted as the coupling of two
707 baroclinic VRWs across the moat. In contrast, this wobbling was not reproduced in a two-
708 dimensional barotropic model framework unless an asymmetric condition is provided at the

709 initial time. The interaction of multiple eyewalls also affects the intensity and other structures
710 substantially.

711 The inner eyewall of Hurricane Wilma (2005) became elliptic (Fig. 12; Lai et al. 2019),
712 and the radial flow had significant wavenumber-2 component at lower levels as
713 wavenumber-1 component weakened. Using a nondivergent barotropic model, they showed
714 that this structure was developed through barotropic instability across the moat in which the
715 sign of radial vorticity gradient changed (referred to as type-2 instability; Kossin et al. 2000).
716 The inner eyewall significantly decays due to the development of wavenumber-2 radial flow.
717 With a three-dimensional full-physics model, Lai et al. (2021b) pointed out that the inner
718 eyewall decay corresponded to the net negative radial advection after the elliptic structure
719 developed through the type-2 instability. It suggests that the eddies that develop under
720 concentric eyewall condition can significantly weaken the inner eyewall, as well as the cutoff
721 of moisture transport in the boundary layer. This finding is supported by the observation-
722 based study of Tsujino et al. (2021).

723 Lai et al. (2021b) also mentioned that an idealized simulation with physics
724 parameterization schemes turned off can intensify the outer eyewall in addition to the inner
725 eyewall decay. Lai et al. (2021c) investigated the interaction between the inner eyewall and
726 outer eyewall using an unforced shallow-water model. The radial eddy transport due to
727 VRWs that developed through the type-2 instability can explain the inner eyewall weakening
728 and outer eyewall intensification. In terms of AAM budget, the inner eyewall weakening and
729 outer eyewall intensification can be viewed as divergence and convergence of the eddy
730 angular momentum flux, respectively.

731 While the above-mentioned studies discuss the early stage of eddies developing
732 according to the type-2 barotropic instability, Lai et al. (2021a) investigated the impacts for
733 a longer period using forced and unforced shallow-water equations. The inner eyewall
734 weakening and outer eyewall intensification repeatedly appear associated with the

735 development of eddies. They also implied that the difficulty in the prediction of intensity
736 changes during ERCs partly stems from type-2 instability.

737 Rostami and Zeitlin (2022) indicated that a robust tripolar vortex structure emerges
738 during the late stages of the type-2 barotropic instability of double eyewall TC-like vortices
739 when the intensities of the eyewalls are comparable with a moist-convective shallow-water
740 model. This tripolar structure is more complicated than those in the studies employing two-
741 dimensional nondivergent barotropic models, rotating shallow water models or laboratory
742 experiments. Slocum et al. (2023b) demonstrated the type-2 instability by using a two-
743 dimensional nonlinear nondivergent barotropic model with a simplified vorticity profile of
744 Hurricane Maria (2017). Their results indicated that the type-2 instability may occur near the
745 end of an ERC and yield a tripolar vorticity structure.

746

747 *3.8 Decay of the inner eyewall*

748 Using tangential winds estimated from Himawari-8 satellite data, Tsujino et al.
749 (2021) demonstrated that momentum loss to surface friction could not sufficiently explain
750 the inner eyewall intensity decay rate during the ERC of Typhoon Trami (2018). They
751 suggested that eddies also weaken the inner eyewall. Based on a series of slab boundary
752 layer model experiments, Kuo et al. (2022) proposed a scaling law in which inner eyewall
753 pumping is proportional to storm intensity and the root square of a non-dimensional moat
754 size (the actual moat size normalized by the Rossby length). From this perspective, they
755 explained why satellite observations often document wider moats in long-lived ERCs (Yang
756 et al. 2021).

757 Using idealized simulations, Yang et al. (2024) showed that the decay of the inner
758 eyewall occurs earlier following the onset of an outer eyewall if the initial vortex is small
759 and weak. They demonstrated that spiral rainbands of a large and strong TC were more
760 active outside of the outer eyewall. These active rainbands reduced the inward-penetrating

761 inflow toward the outer eyewall, slowing down the contraction and intensification of the
762 outer eyewall. This prolonged the duration of the concentric eyewall structure.

763

764 *3.9 Special ERCs cases*

765 Molinari et al. (2019) summarized several unusual features of ERCs during
766 Hurricane Frances (2004). It underwent three consecutive ERCs that resulted in an
767 annular eyewall structure, which is rarely observed during Pacific typhoons. It has been
768 proposed that unusual easterlies in the upper levels produce an environment with low
769 inertial stability. This persistently low inertial stability in the outflow layer contributes to an
770 anomalously strong outflow and faster intensification of the outer eyewall, resulting in
771 multiple ERCs and shorter ERC durations.

772 Using airborne radar observations, Razin and Bell (2021) analyzed the
773 unconventional ERC of Hurricane Ophelia (2005), which underwent two ERC events at
774 category-1 intensity and over a low SST, at approximately 23°C. Their results supported two
775 previously proposed pathways to SEF; one is due to stratiform heating of TC rainbands
776 (Didlake et al. 2018), while the second is associated with the unbalanced boundary layer
777 dynamics (Huang et al. 2012).

778 Wang and Wang (2021) conducted sensitivity experiments on the decay of the
779 original eyewall and the genesis of an outer eyewall as TC Megi (2011) crossed Luzon Island.
780 The original inner eyewall was not fully destroyed over flat land, highlighting the importance
781 of topography. Additionally, short-lived deep convection (possibly a new inner eyewall)
782 developed after the center of TC Megi (2011) moved into the South China Sea due to
783 favorable oceanic conditions for deep convection. Lau et al. (2024) also investigated island-
784 induced eyewall formation in the case of TC Mangkhut (2018). Their findings generally align
785 with those of Wang and Wang (2021). Furthermore, Lau et al. (2024) analyzed their results
786 from the perspective of unbalanced dynamics.

787

788 *3.10 Other perspectives*

789 The impact of TC-induced SST cooling, ocean waves, radiation, and data
790 assimilation on the SEF and ERC has been addressed in several studies. Yang et al.
791 (2020) and Li et al. (2022a) demonstrated that the inclusion of TC induced SST cooling
792 and ocean waves led to a simulated ERC duration that better matched observations. Yang
793 et al. (2020) displayed that when both the ocean and waves were considered, the lifetime
794 of the ERC was significantly prolonged for Typhoon Sinlaku (2008). According to Li et al.
795 (2022a), the exclusion of the SST cooling associated with slow TC motion resulted in an
796 impractically long ERC during Typhoon Trami (2018). In this case, the substantial negative
797 feedback likely reduced the energy supply to the inner eyewall more severely. Because
798 these results seem oppositional, further investigation is needed to clarify the impact of a
799 coupling as well as the relevant mechanism.

800 Trabing and Bell (2021) showed that the magnitude of shortwave radiation affects
801 the SEF timing through nonlinear interactions at longer lead times. Generally, shortwave
802 radiation served to delay the SEF and ERC. The proposed mechanism is that the
803 shortwave radiation alters the convective and stratiform heating profile and frequency, and
804 it stabilize the atmosphere and reduces the convective available potential energy.

805

806 **4 Operational Analysis and Forecast**

807 Utilizing information about the inner core structure such as RMW in an operational
808 setting to forecast structural changes and the wider implications for TC intensity prediction
809 is challenging. Operational centers have assessed inner core structural changes with the
810 aid of existing and developing operational tools. Combining subjective assessments of
811 inner-core changes with the available objective guidance remains the leading operational
812 methodology. Although all available data are useful for the assessment by forecasters and

813 analysts, the availability and manual assessment of satellite imagery is essential. The
814 inner core structural change is important for the TC intensity estimate because satellite
815 analysts at operational centers examine satellite images and utilize the Dvorak Technique
816 to derive a current intensity (CI) number, which is one of important sources of TC intensity
817 (Dvorak 1984). In this section, the operational analysis and forecasting of the TC inner
818 core structure, use of tools, and relevant research are described. Regarding the surface
819 wind estimation, readers may refer to Knaff et al. (2021), which is a nice review for
820 operational center practices, historical databases, current and emerging objective
821 estimates of TC surface winds, including algorithms, archive datasets, and individual
822 algorithm strengths and weaknesses.

823

824 *4.1 Operational analysis of RMW*

825 The RMW can be operationally used for various purposes: 1) to monitor inner core
826 structural changes associated with SEF, ERC, and RI; 2) to estimate central pressure; 3)
827 to correct numerically simulated wind fields (e.g., Aijaz et al. 2019), 4) to potentially aid
828 diagnostic intensity forecasts (e.g., Carrasco et al. 2014), and 5) to forecast waves and
829 storm surges. Considering the importance of the RMW to the TC wind structure, a
830 questionnaire was conducted to collect operational practices on the method of RMW
831 analyses. The questions were:

832 *“Q1: Does your operational center analyze the RMW routinely and include it in your post-
833 analysis best track?”*

834 *Q2-1: If yes to Q1, please describe the method of analysis with possible references such
835 as research papers, documents, memorandum, etc.*

836 *Q2-2: If no, please select the reasons (check all that apply). 1. Lack of observations, 2.
837 Lack of methods, 3. No needs, 4. Others”*

838 The respondents to the questionnaire were Regional Specialized Meteorological

839 Center (RSMC) Miami, RSMC Tokyo, RSMC La Réunion, RSMC Nadi, Tropical Cyclone
840 Warning Centre (TCWC)-Australia, TCWC-Wellington, and Joint Typhoon Warning Center
841 (JTWC). Of the seven respondents, four indicate that their center analyzes RMW in real-
842 time, two centers do not analyze RMW, and one center analyzed RMW in real time
843 however they do not conduct post-analysis/quality control. The answers to these questions
844 are summarized in Tables 1 and 2.

845 For centers that generally analyze RMW, a manual assessment of all available
846 observations was the primary method. Radar, scatterometry, passive microwave (37 GHz
847 and 89 GHz frequencies), and visible imagery, as well as Synthetic Aperture Radar (SAR),
848 Soil Moisture Active/Passive (SMAP), and Soil Moisture and Ocean Salinity Mission
849 (SMOS)-derived wind speeds are analyzed, with preference for the highest quality data.
850 Climatology is occasionally used in the absence of real-time observations. Internal
851 documents describe the best practices for analysis, particularly for microwave imagery, as
852 well as the procedures for shifting from the inner RMW to the outer RMW as the primary
853 radius of maximum winds when an ERC is taking place. Post-analysis of the RMW is
854 performed coincidentally with other parameters, such as position and intensity. Centers
855 that do not routinely analyze RMW cited the sparsity of observations, lack of operational
856 requirements, non-standard structures such as monsoon gyres, and temporal
857 inhomogeneity of supporting datasets.

858 From a research community, Tsukada and Horinouchi (2023) recently proposed
859 the improved method of IR-based RMW estimation for a TC with a clear eye, following up
860 the work of Kossin et al. (2007). Compared with C-band SAR sea-surface wind estimates,
861 the mean absolute error of 4.7 km in previous studies was reduced to 1.7 km. Chavas and
862 Knaff (2022) developed a simple semiempirical model to estimate the RMW using
863 operationally available parameters, including an outer wind radius, the Coriolis parameter,
864 and maximum wind speed. The empirical model is based on the physical understanding of

865 the radial structure of TCs (Chavas and Lin 2016; Emanuel and Rotunno 2011; Emanuel
866 et al. 2004). Chavas and Knaff (2022) showed that the model estimates RMWs with much
867 better accuracy than previous methods (e.g., Knaff et al. 2011; Knaff et al. 2015).
868 Additionally, Avenas et al. (2023) applied the method of Chavas and Knaff (2022) to
869 estimate RMWs using outer wind radii derived from satellite radiometers and
870 scatterometers and SAR-derived RMWs as ground truth.

871

872 *4.2 Operational assessment of ERCs*

873 Although the predictability of an ERC event is low, efforts have been continuing.
874 As ERC affects intensity trends and anticipated impacts, predicting the onset and
875 completion of ERC remains an important challenge for operational centers. While high
876 resolution numerical weather prediction models can resolve ERCs, the variability in
877 predictions and low skill limit their applicability in operational settings. In many cases, the
878 models indicated an ERC event, however the timing is incorrect. Applying hourly data
879 assimilation cycling to the simulations of Hurricane Matthew (2016), Green et al. (2022)
880 showed that concentric eyewalls could be better resolved by the inclusion of ground-based
881 radar observations. Assimilating radar observations more rapidly reduced the bias of the
882 storm structure, indicating the importance of an improved representation of the initial TC
883 structure in forecasting the ERC process.

884 Considering the limitations of high-resolution modeling, forecasters at National
885 Hurricane Center (NHC), Central Pacific Hurricane Center (CPHC), JTWC, RSMC La
886 Réunion and Bureau of Meteorology (BoM) assess the University of Wisconsin CIMSS
887 Microwave Probability of Eyewall Replacement Cycle (M-PERC) tool (Kossin et al. 2023;
888 available online at https://tropic.ssec.wisc.edu/real-time/archerOnline/web/index_erc.shtml)
889 to assist ERC prediction. M-PERC, which was developed using Atlantic data, applies a
890 logistic regression probabilistic model to evaluate the environmental conditions that favor

891 ERC. Evaluation of M-PERC indicate that the algorithm efficiently depicts ERC events in
892 all ocean basins despite the Atlantic focused training of the model. Independent validation
893 of the model using a climatological probability of ERC of 13% yielded a Brier skill score of
894 35%. Additionally, Pulmano and Joykuty (2021) found that for Atlantic Basin TCs between
895 2017 and 2019, the algorithm correctly predicted approximately 41% of the total ERC
896 events. The model was trained on SEF events that lead to ERC and thus model
897 probabilities will rise with SEF development and intensification of the outer ring regardless
898 of where the ERC completes or not. Pulmano and Joykuty showed that this can at times
899 lead to higher probability events in the presence of dry air environments at higher latitudes
900 that do not result in completion of the ERC. In addition, Kossin et al. (2023) noted the
901 intensity evolution of ERC events as a function of TC intensity and found that the previous
902 paradigm of TC weakening as result of an ERC should be modified to note a change in
903 intensification rate instead. Kossin et al. (2023) shows TC intensification rate changes
904 during ERC events which indicates a slowing of intensification for storms undergoing
905 ERCs for Category 1-2 storms whereas storms at Category 3 intensity or stronger are
906 more likely to undergo weakening (Fig. 13). Lower probability ERC events indicated by M-
907 PERC may occur during the RI phase with the ERC resolving quickly and only resulting in
908 a brief pause in RI on the order of 12 hours with RI resuming once the ERC is completed.
909 An example of this occurred during the RI of Hurricane Dorian in 2019 on August 29th
910 which was confirmed by aircraft reporting concentric eyewalls with the inner eye diameter
911 of only 4 nautical miles.

912

913 *4.3 Use of inner core structure data for operational intensity prediction*

914 Operational centers either use or anticipate tools such as M-PERC, raw
915 microwave satellite imagery, statistical-dynamical models, and lightning data to predict the
916 intensity fluctuations associated with changes in the inner core structure. A case study in

917 the 2020 JTWC Annual Tropical Cyclone Report highlighted the successful prediction of
918 multiple ERC events for TC 25P (Harold) using M-PERC algorithms, demonstrating the
919 value of the tool for anticipating the associated short-term intensity fluctuations (typically a
920 slight weakening prior to re-intensification) in an operational setting (Francis and Strahl
921 2021). Additionally, the BoM reported frequent use of microwave imagery to provide short-
922 term subjective predictions of inner core structural changes and related intensity
923 fluctuations. For example, a visual assessment of the ERC evident in passive microwave
924 imagery influenced the operational decision-making for TC Trevor (2019). The system was
925 approximately 12 h from landfall at CI 5.0 / 80 knots in a very favorable broad scale
926 environment. Intensification at or above the standard Dvorak rate would lead to Category 4
927 (Australian system – 90-105 knot 10 min. mean wind) landfall forecasts were credible.
928 However, based on the double eyewall structure evident in the SSMI microwave imagery
929 (Fig. 14), a subjective judgment was made that the inner eyewall would likely decay over
930 the next 12-18 h leading to little intensification over that period, therefore the landfall
931 forecast was limited to Category 3. In this case, this is what unfolds (Tony Wedd, Senior
932 Meteorologist, BoM).

933 Rapid intensification can be preceded by inner core changes, as shown by Li et al.
934 (2022b) who used the best-track dataset for the North Atlantic and Eastern North Pacific
935 during 1999-2019 to examine the statistical relationship between the rapid contraction of
936 the RMW and RI. Their findings demonstrated that rapid RMW contractions were
937 frequently followed by RI. Operationally, forecasters subjectively assess passive
938 microwave imagery and other satellite imagery to identify processes such as contraction of
939 the inner core as indicators of intensity change. Intensity tools aim to capture these
940 processes by placing various levels of emphasis on each process, depending on their
941 predictive value. The well-known statistical hurricane intensity prediction scheme (SHIPS;
942 DeMaria et al. 2005) incorporated the IR brightness temperature parameter derived from

943 infrared satellite data. The Rapid Intensification Prediction Aid (RIPA), which provides
944 statistical guidance for predicting the likelihood of RI in the western North Pacific and other
945 basins, incorporates IR brightness temperatures, core size, and core symmetry as
946 predictor variables. RIPA has been used operationally at the JTWC since late 2017 (Knaff
947 et al. 2020).

948 Observational data have been used to validate physical parameterizations in
949 operational TC forecast models. Composites of global positioning system (GPS)
950 dropsondes, Doppler radar and Stepped Frequency Microwave Radiometer (SFMR) data
951 have been used to evaluate the performance of the forecasted TC structures which lead to
952 identification of model deficiencies (e.g. Zhang et al. 2015). Observation based new
953 parameterizations of turbulent mixing have been implemented in U.S. operational TC
954 models (e.g., Hurricane Weather and Research Forecasting System (HWRF)), which led
955 to significant improvement in hurricane intensity forecasts as well as RI prediction (Zhang
956 et al. 2015; Zhang et al. 2018).

957 Operational forecasters consider the potential applicability of inner core lightning
958 data for their forecast processes, particularly considering the limited availability of high-
959 resolution microwave imagery, which is typically necessary to reveal rapidly changing
960 inner core structures. The lightning data from TC Harold (2020), which underwent multiple
961 periods of rapid intensification, indicated patterns that were highly consistent with the
962 documented relationships between increased inner core lightning density and
963 intensification (Francis and Strahl 2021; Lin and Chou 2020; Stevenson et al. 2018).

964 Operational forecasters consider lightning density as a possible forecasting aid, and the
965 lightning-based model is running in parallel and demonstration mode at NHC for the
966 Atlantic and eastern Pacific using GOES-16 and GOES-18 GLM data (Slocum et al.
967 2023a). Further studies are needed to identify reliable lightning indicators that are valid for
968 all basins and can be used as tools to predict intensity changes, particularly for RI.

969 Finally, the ingestion of dry air into the inner core is associated with rapid
970 weakening. Forecasters report using satellite imagery, including Morphed Integrated
971 Microwave Imagery at the Cooperative Institute for Meteorological Satellite Studies
972 (MIMIC) total precipitable water loops, to identify inner core erosion due to dry air intrusion
973 and VWS. However, assessing satellite imagery is a subjective process, and reliably
974 determining the extent and effects of dry air intrusion into the inner core remains a
975 forecasting challenge.

976

977 **5 TC inner core structure change in future climate**

978 The horizontal scale of a TC affects the size of the damaged area caused by
979 strong winds and waves, attracting socioeconomic interest. Future changes in TC inner
980 core structure are also an important issue. Some studies have investigated future changes
981 in TC size (Knutson et al. 2020). Gutmann et al. (2018) evaluated the influence of global
982 warming on the size of TCs using the WRF model and indicated that although the
983 influence depended on individual TCs, the mean size across all TCs did not change
984 significantly owing to global warming. Wehner (2021) reported that climate change does
985 not influence the average radial distribution of a simulated TC on a specified Saffir-
986 Simpson scale, at least when using the NCAR Community Atmosphere Model version 5.1
987 with a horizontal grid spacing of 25 km. Song et al. (2020) showed that TC size would
988 increase under quadruple CO₂ forcing over all ocean basins except the North Atlantic and
989 North Indian Oceans using SEM0-UNICOM. As introduced here, the results obtained in
990 previous studies are inconsistent, which is partly due to differences in the definition of TC
991 size, experimental design, and model. Knutson et al. (2020) emphasized that “future
992 studies should further assess model capabilities of simulating present day TC sizes, which
993 has so far been done only to a limited extent. A better understanding of the mechanisms
994 determining TC sizes in observations and models is important, as is the monitoring and

995 accumulation of observed climate records of TC size.”

996 In general, TCs are in a mature stage in the tropics and consist of a strong rotating
997 circulation with a large axisymmetric component. As TCs transition into the mid-latitudes,
998 their structures become more asymmetric owing to the baroclinic environment in the
999 midlatitudes. Typhoon Faxai in 2019 made landfall in the Tokyo metropolitan area at a
1000 latitude of approximately 35°N with a central pressure of 960 hPa, and caused severe
1001 damage owing to strong winds. The intensity of Faxai exceeded the maximum potential
1002 intensity, which was attributed to Faxai maintaining an axisymmetric structure similar to
1003 that of a well-developed TC in the tropics well after a TC would normally begin an
1004 extratropical transition. This was due to favorable environmental conditions such as
1005 relatively weak VWS (Miyamoto et al. 2022). It is unclear whether the environmental
1006 conditions before Faxai made landfall in Japan are associated with global warming.
1007 However, there is an urgent need to address the impacts of global warming on TC
1008 structure in the mid-latitudes where large cities and densely populated areas are located.
1009 Using a nonhydrostatic regional model with a horizontal resolution of 0.04°, Kanada et al.
1010 (2020) conducted approximately 100 dynamical downscaling experiments for mid-latitude
1011 typhoons in both current and warming climates (Fig. 15). In a future warmer climate, the
1012 extratropical transition position of tropical cyclones will shift to higher latitudes owing to
1013 higher SST, larger near-surface water vapor content/capacity, and smaller baroclinicity
1014 compared to the current climate conditions. The results of Fujiwara et al. (2023) were
1015 consistent with these results; they also showed that eyewalls become deeper in warmer
1016 climates. This will likely facilitate TCs intruding further into mid-latitude regions while
1017 retaining an axisymmetric structure, intensity, and smaller radius of maximum winds, as
1018 observed in the tropics.

1019 Changes in several recent destructive TCs have also been studied using
1020 convection-permitting regional climate model simulations under pre-industrial, current, and

1021 future climate conditions (Patricola and Wehner 2018). The results indicate significant
1022 rainfall contraction and enhancement in the core regions of storms simulated in future
1023 warmer climate conditions and suggest that climate change has likely begun to enhance
1024 rainfall for recent destructive TCs.

1025

1026 **6. Concluding remarks**

1027 This paper compiled recent publications and relevant information on research and
1028 operational use of TC inner core structural changes with an overview of conventional
1029 understanding. These are summarized as follows:

- 1030 ● The changes in the inner core structure during each intensity stage are described in
1031 detail. Many studies have focused on non-conventional distinctive structures such as
1032 dual warm cores, mid-tropospheric maxima in wind speed, short-lived deep convective
1033 clouds in the eye region, and finger-like features from both observational and modeling
1034 perspectives. The numerical simulations confirm the dependence of the RMW on the
1035 size and humidity of the outer core region. Non-axisymmetric structures, including
1036 those related to VWS, as well as contributions from factors such as translation speed,
1037 interactions with the land and ocean, and transformations from high-frequency
1038 components have advanced. A comparison with dry TCs suggests the potential impact
1039 of falling hydrometeors on the depth of the inflow layer and the size of the inner core.
- 1040 ● Observations and numerical simulations have focused on the transformation of the
1041 SBC into a secondary eyewall. Diabatic cooling from rain in the SBC causes MDI,
1042 which triggers enhanced convection in the upshear region. The axisymmetrization of
1043 the enhanced convection finally yielded a secondary eyewall. Also, many studies have
1044 reconfirmed the importance of boundary layer dynamics in the SEF region. From a
1045 dynamical perspective, the interaction between the inner and outer eyewall vortices
1046 has been applied to explain the enhanced wobbling and decay of the inner eyewall or

1047 the persistence of multiple eyewalls. Some studies have investigated the
1048 environmental conditions required for special multiple eyewall events. As such, both
1049 dynamical and thermodynamical processes have been proposed in the boundary layer
1050 and free atmosphere for multiple eyewalls, and the understanding of the SEF and
1051 ERC has progressed over the last several years. However, further studies are needed
1052 to clarify whether they are complementary or exclusive, and whether they are case
1053 dependent.

1054 ● With the progress in understanding the inner core structure, many operational
1055 agencies have started to utilize information on the inner core structure for disaster
1056 prevention and mitigation. The questionnaire to operational centers revealed that five
1057 operational agencies had already analyzed the RMW of TCs based on available
1058 observations with a preference for the highest quality data. Some agencies have
1059 begun to use guidance and subjective analyses to analyze and/or predict SEF and
1060 ERC. Inner core lightning activity and dry air intrusion from satellite observations were
1061 investigated for use in operational analysis and forecast. In-situ observations were
1062 used to improve model physics with a focus on turbulent mixing.

1063 ● Studies on climatological changes in the inner core structure are limited or unreliable
1064 owing to the high computational cost and diversity of the results. However, several
1065 publications have reported an increase in the inner core rainfall amount, an
1066 axisymmetric structure maintained at higher latitudes using regional climate models.

1067
1068 Building upon prior studies, extensive research has been conducted on structural changes
1069 in the TC inner core during various phases, SEF and ERC. Consequently, our knowledge
1070 of TC inner core structure has significantly advanced in recent years. However, some
1071 controversies and fragmented understandings still remain. It is essential to appropriately
1072 assess each result in terms of methodological reliability and strive for a more generalized

1073 perspective. It requires refinement in observational methodologies and numerical
1074 simulations, as the TC inner core is not easily observable. On one side, some operational
1075 agencies have already begun analyzing the RMW and concentric eyewalls using both
1076 existing and new tools. Although standardizing methodologies is challenging, enhanced
1077 international communication among operational centers and researchers is expected to
1078 foster advancements in publishing valuable information for TC-related disaster prevention
1079 based on updated insights. Further research on the impact of climate change on TC inner
1080 core structure is recommended. To make this more feasible, further model improvements
1081 and development of a kilometer-scale climate model (Miura et al. 2023) will be required,
1082 which demands huge computational resources. It is also worthy of note that machine
1083 learning and artificial intelligence techniques have been growing rapidly as a tool which
1084 may yield the new insight on the TC inner core research and operation.
1085

1087 Abbreviations

- 1088 AAM – Absolute Angular Momentum
- 1089 AMSR2 – Advanced Microwave Scanning Radiometer-2
- 1090 ARCHER – Automated Rotational Center Hurricane Eye Retrieval
- 1091 ASCAT – Advanced SCATterometer
- 1092 BoM – Australian Bureau of Meteorology
- 1093 CI – Current Intensity
- 1094 CPH – Central Pacific Hurricane Center
- 1095 CIMSS – Cooperative Institute for Meteorological Satellite Studies
- 1096 DA – Data Assimilation
- 1097 ERC – Eyewall Replacement Cycle
- 1098 GPS – Global Positioning System
- 1099 HWRF – Hurricane Weather Research and Forecasting
- 1100 IR – InfraRed
- 1101 JTWC – Joint Typhoon Warning Center
- 1102 LES – Large Eddy Simulation
- 1103 MDI – Mesoscale Descending Inflow
- 1104 MIMIC – Morphed Integrated Microwave Imagery at the Cooperative Institute for
1105 Meteorological Satellite Studies
- 1106 M-PERC – Microwave Probability of Eyewall Replacement Cycle
- 1107 NHC – National Hurricane Center
- 1108 NRL – Naval Research Laboratory
- 1109 RI – Rapid Intensification
- 1110 RIPA – Rapid Intensification Prediction Aid
- 1111 RMW – Radius of the Maximum Wind
- 1112 RSMC – Regional Specialized Meteorological Center

- 1113 SAR – Synthetic Aperture Radar
- 1114 SBC – Stationary Band Complex
- 1115 SEF – Secondary Eyewall Formation
- 1116 SFMR – Stepped Frequency Microwave Radiometer
- 1117 SHIPS – Statistical Hurricane Intensity Prediction Scheme
- 1118 SMAP – Soil Moisture Active/Passive
- 1119 SMOS – Soil Moisture and Ocean Salinity mission
- 1120 SST – Sea Surface Temperature
- 1121 TC – Tropical Cyclone
- 1122 TCWC – Tropical Cyclone Warning Centre
- 1123 T-PARCII – Tropical Cyclones–Pacific Asian Research Campaign for the Improvement of
1124 Intensity estimations/forecasts
- 1125 VRW – Vortex Rossby Wave
- 1126 VWS – Vertical Wind Shear (i.e., Environmental vertical shear of horizontal wind)
- 1127

1128 **Data availability statement**

1129 This paper is a review article based on the recent publications, answers to questionnaires,
1130 and experiences of forecasters. See the paper and/or contact the corresponding author of
1131 the paper for the detailed data availability. Regarding the questionnaires and forecasters'
1132 experiences, contact Kosuke Ito for more details.

1133

1134

1135 Acknowledgement

1136 This paper was based on the report “TC structure change processes: Inner core” for 10th
1137 International Workshop on Tropical Cyclones (IWTC-10). We appreciate the contribution of
1138 Dr. Margie Kieper, who passed away on 25 April 2024. Kosuke Ito was supported by JST
1139 Moonshot R&D Grant Number JPMJMS2282-06 and JSPS KAKENHI Grant Numbers
1140 JP21H04992 and JP23K26359. Yoshiaki Miyamoto was supported by JSPS KAKENHI
1141 Grant Numbers JP18H05872, JP19H05696, JP19H01973, JP19K04849, and
1142 JP19K24677. Chun-Chieh Wu and Yi-Hsuan Huang was supported by the National
1143 Science and Technology Council of Taiwan through Grant NSTC 112-2123-M-002-002,
1144 and Chun-Chieh Wu was also supported by the Office of Naval Research through Grant
1145 N00014-20-1-2467. James Hlywiak was supported by U.S. National Science Foundation
1146 (NSF) PREEVENTS Track 2 Award 1663947. Yohei Yamada was supported by JST
1147 Moonshot R&D Grant Number JPMJMS2282-10, JSPS KAKENHI Grant Numbers
1148 JP20H05728, and MEXT as “Program for promoting researches on the supercomputer
1149 Fugaku” JPMXP1020200305. Jun Zhang was supported by U.S. NOAA Grants
1150 NA21OAR4590370, NA22OAR4590178, and NA22OAR4050669D, Office of Naval
1151 Research Grant N00014-20-1-2071, and U.S. NSF Awards 2228299 and 2211308. Sachie
1152 Kanada was supported by JSPS KAKENHI Grant Numbers JP19H05696, JP20H05166
1153 and MEXT-Program for the advanced studies of climate change projection SENTAN Grant
1154 Number JPMXD0722680734. We thank to Dr. Christopher Slocum who gave us comments
1155 on the operational use of lightning-based guidance.

1156

1157
 1158 Table 1. The responses from operational centers to Q2-1 in Section 4.1.
 1159

Response #	Explanations
1	Manual analysis of imagery from scatterometers, SMAP, SMOS, SAR, radar, and microwave overpasses
2	All available observations. By decreasing order of quality: radar data, SAR data, scatterometer data, microwave imagery, IR or Vis imagery. Internal documents describe the best way to use microwave imagery in particular, or when to shift from inner RMW to outer RMW as the primary radius of max winds when an ERC is taking place.
3	Standard Dvorak techniques and available scatterometry observations. Otherwise lack of observations.
4	Available radar and microwave imagery. The lower 37GHz channel is preferred, otherwise estimate distance inside the 89 GHz circulation. IR and Vis imagery are used if an eye is present. Additionally, analysis of guidance from scatterometry, noting the various sensors and resolution limitations, each sensor is considered on merit. SMAP, SMOS, ASCAT, AMSR2 and HY-2B can all be helpful, particularly when the RMW is larger. Finally in the absence of objective information, climatology can be used.
5	The post analysis of RMW is performed at the same time as other post analysis of intensity, track, and wind radii using all available observations, etc.

1160
 1161

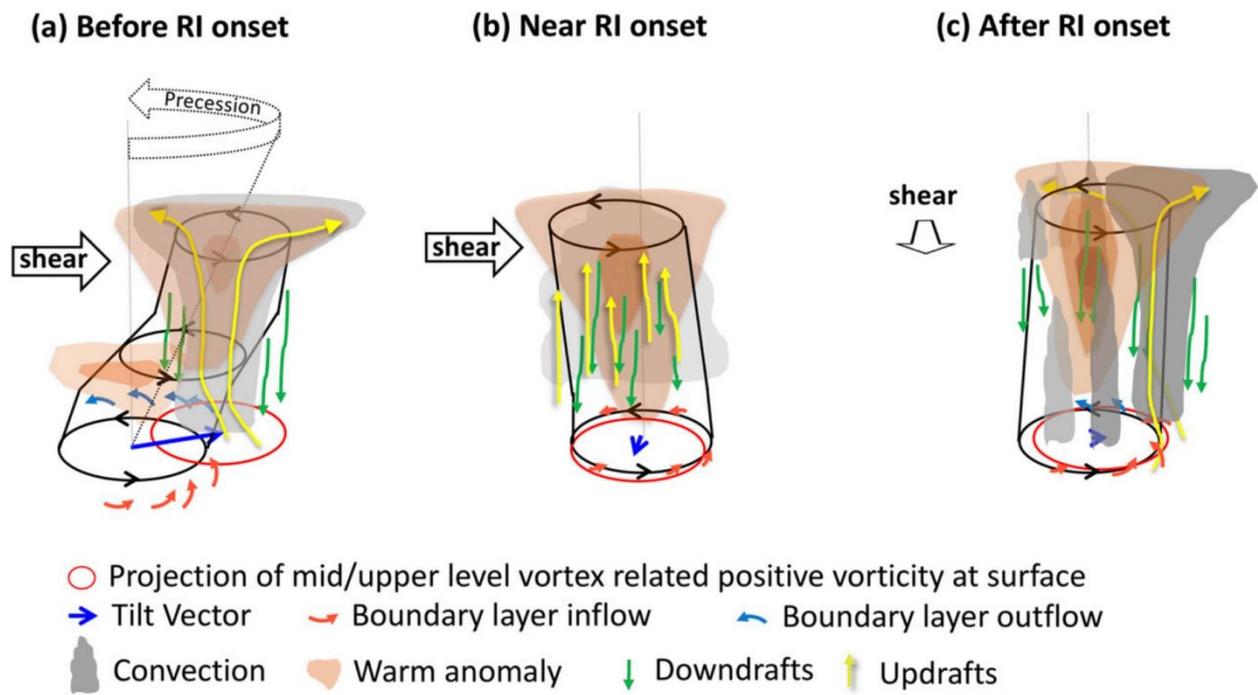
1162 Table 2. The responses from operational centers to Q2-2 in Section 4.1.

1163

Response #	Explanations
1	Lack of observations and no needs
2	While we understand the usefulness of RMW in surveys and research, we believe that careful consideration is needed in terms of the characteristic of the basin such as an unorganized tropical cyclone from a monsoon gyre, and temporal homogeneity of data in order to routinely include RMW in the best track.

1164

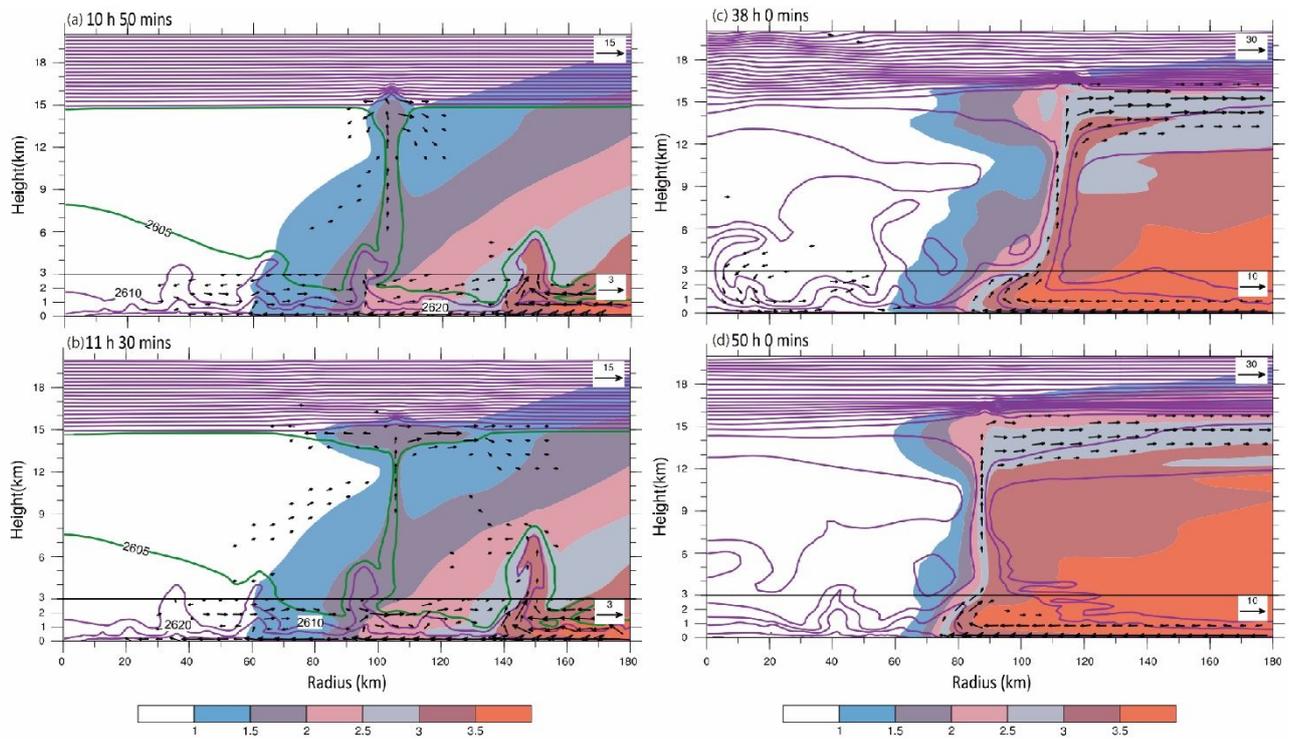
1165



1166

1167 Fig. 1. The adjustment of vortex column, warming anomalies and updrafts during the
 1168 transition period from before RI onset to after RI onset. Figure 16 of Tao and Zhang (2019).

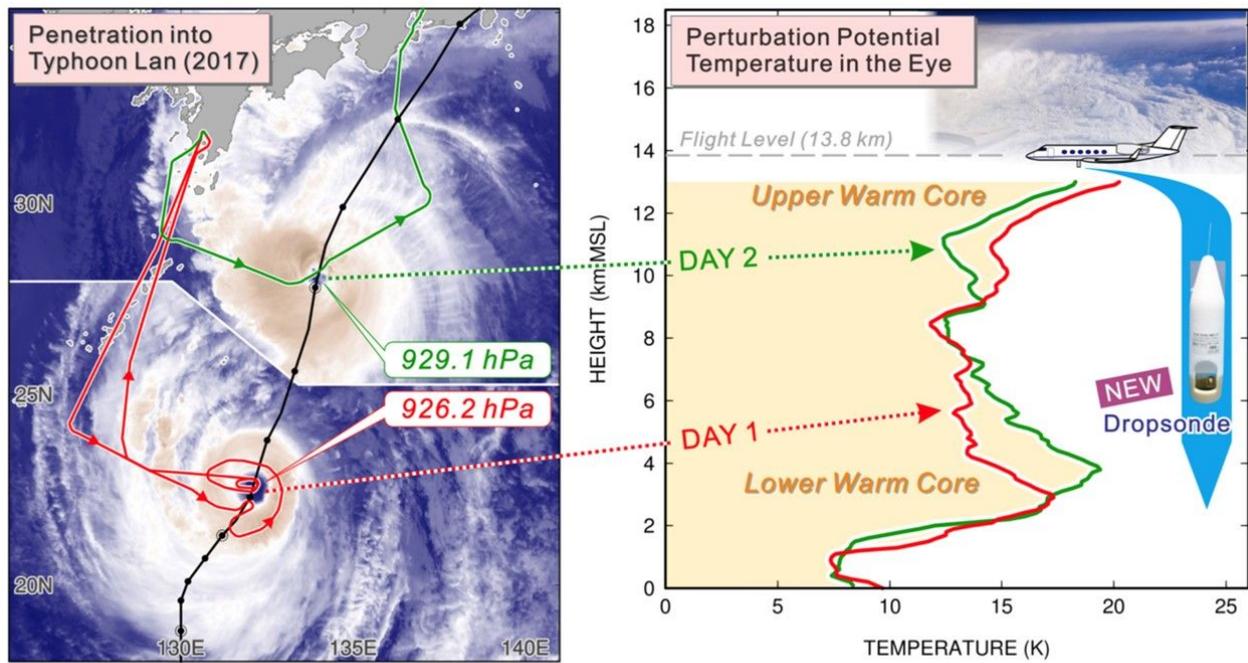
1169 © American Meteorological Society. Used with permission.



1171

1172 Fig. 2. Moist entropy (contour interval = $10 \text{ J kg}^{-1} \text{ K}^{-1}$; purple lines) and angular momentum
 1173 M ($10^6 \text{ m}^2 \text{ s}^{-1}$; shading) at different simulation times. Left column for early stage (Phase I),
 1174 right column for late stage (Phase II). Figures 5a, 5c, 8a and 8c of Peng et al. (2019) with
 1175 labels modified. © American Meteorological Society. Used with permission.

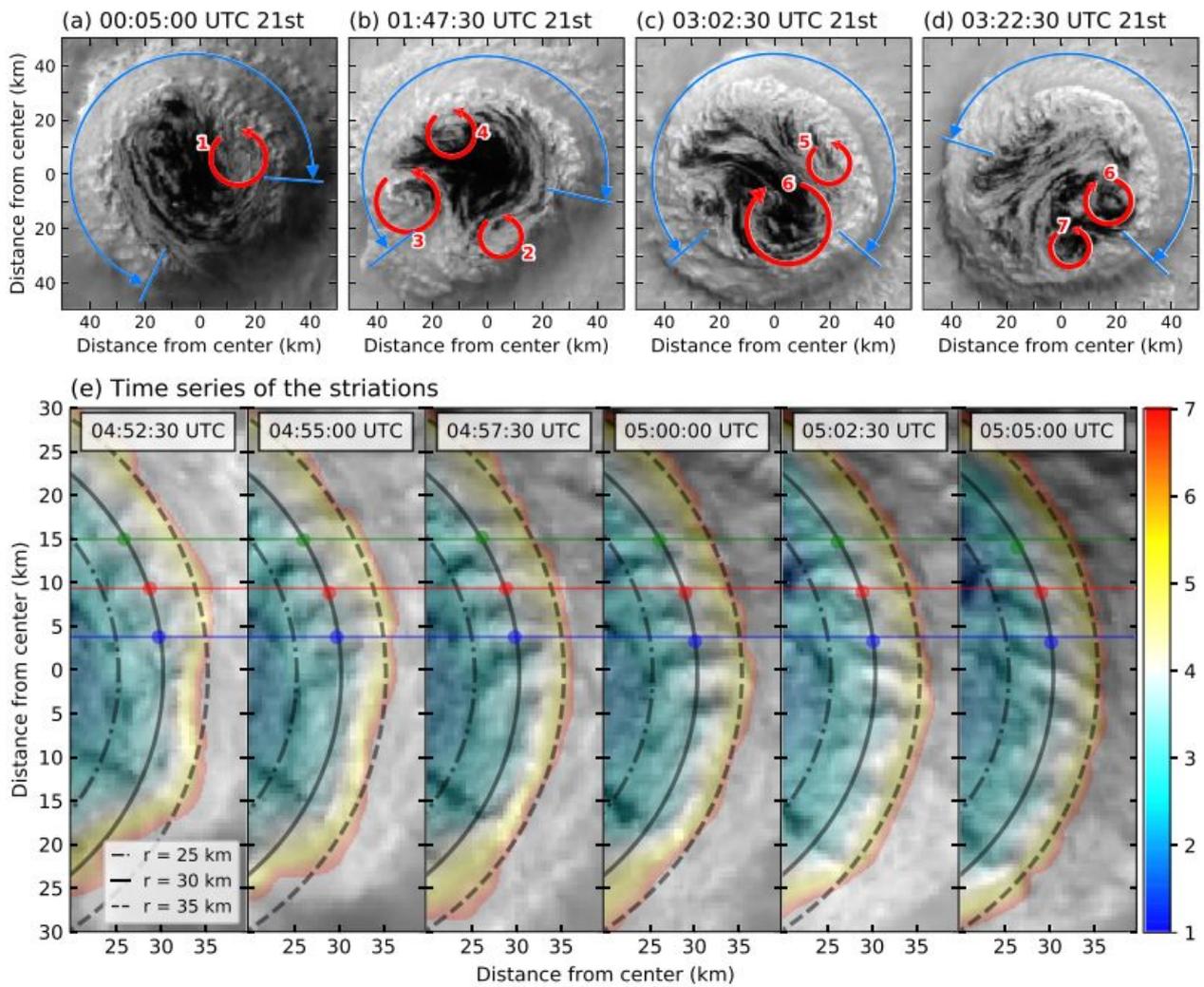
1176



1177

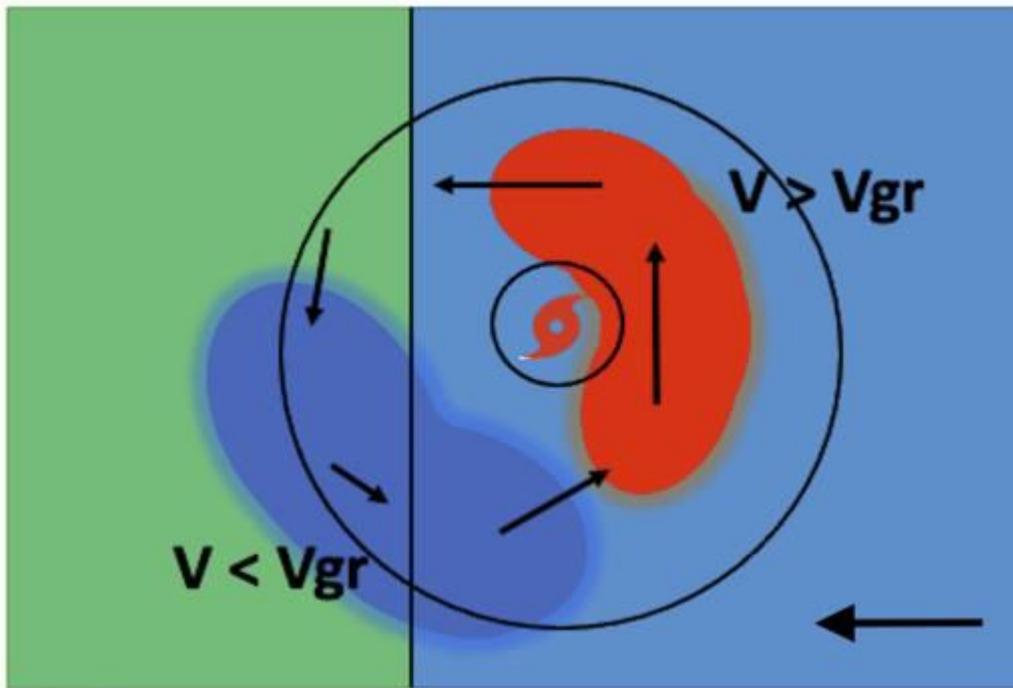
1178 Fig. 3. (left) Flight paths superimposed onto infrared images during the aircraft observation.

1179 (right) Vertical profiles of the observed perturbation potential temperature within the eye.



1181

1182 Fig. 4. Mesovortices and striations. (a-d) Red arrows indicate the mesovortices, and the
 1183 range of striations were indicated by blue arrows. (e) Striations. The colored horizontal lines
 1184 are shown to indicate their positions at the same r if they are rotated at 1.75×10^{-3} rad/s.
 1185 Figure 5 of Tsukada and Horinouchi (2020).



1187

1188 Fig. 5. A schematic illustration of the flow around a landfalling TC. The subgradient winds
1189 with large inflow angles are indicated in blue, while the supergradient winds at the top of the
1190 boundary layer are indicated in red. Figure 23 of Hlywiak & Nolan (2022). © American
1191 Meteorological Society. Used with permission.

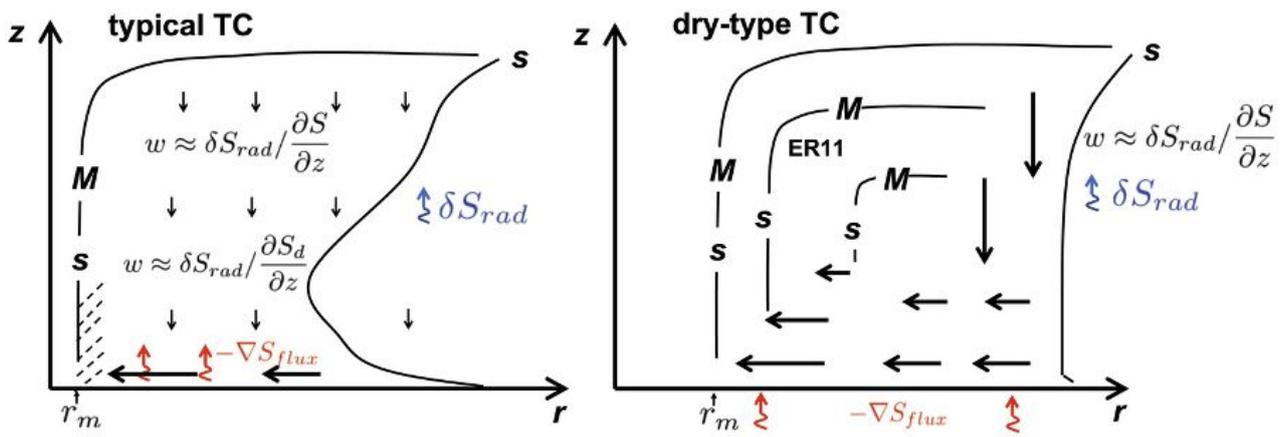
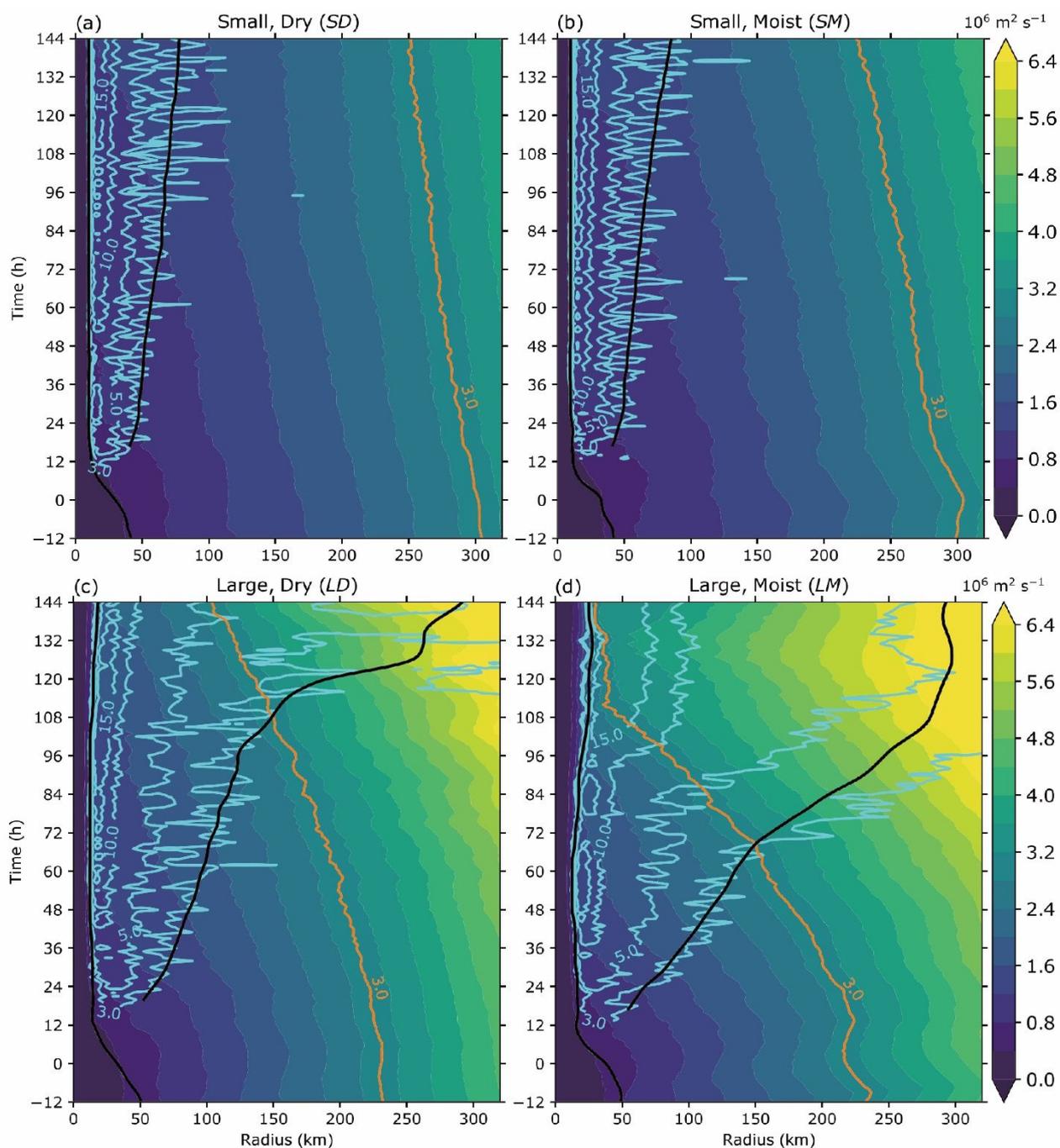
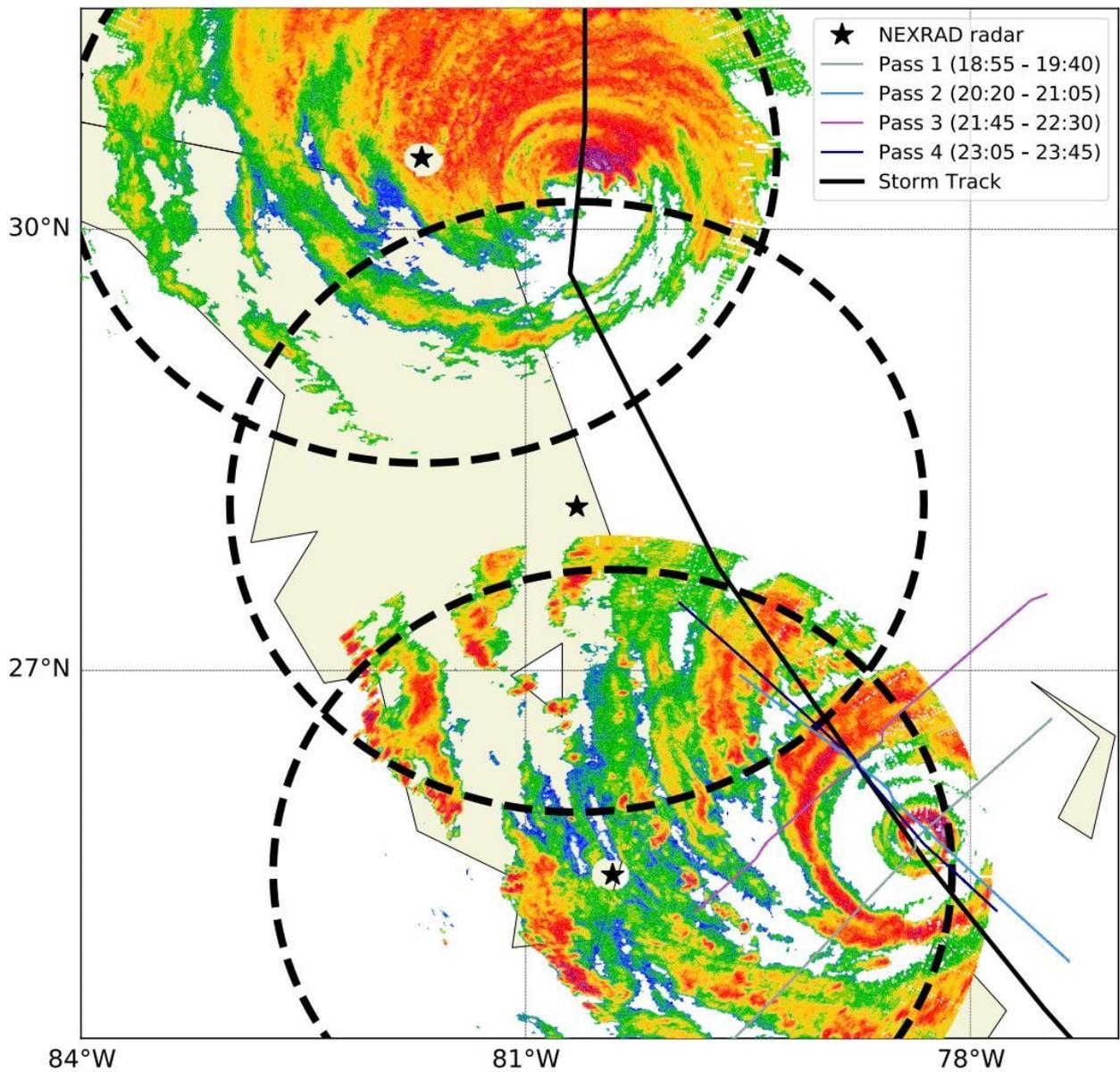


Fig. 6. Schematic plot of typical and dry-type TCs. Figure 16 of Wang and Lin (2020). © American Meteorological Society. Used with permission.



1197

1198 Fig. 7. AAM at $z = 1$ km (shaded), radial velocity averaged from surface to $z = 1$ km (cyan
 1199 contours at 3, 5, 10, and 15 m/s) for the simulations: (a) small incipient vortex in the dry
 1200 environment, (b) small incipient vortex in the moist environment, (c) large incipient vortex in
 1201 the dry environment, and (d) large incipient vortex in the moist environment. The AAM
 1202 contour of $3.0 \times 10^6 \text{ m}^2 \text{ s}^{-1}$ is shown in orange. The inner black line represents RMW, while
 1203 the outer black line represents the radius of gale-force winds. Figure 6 of Martinez et al.
 1204 (2020).



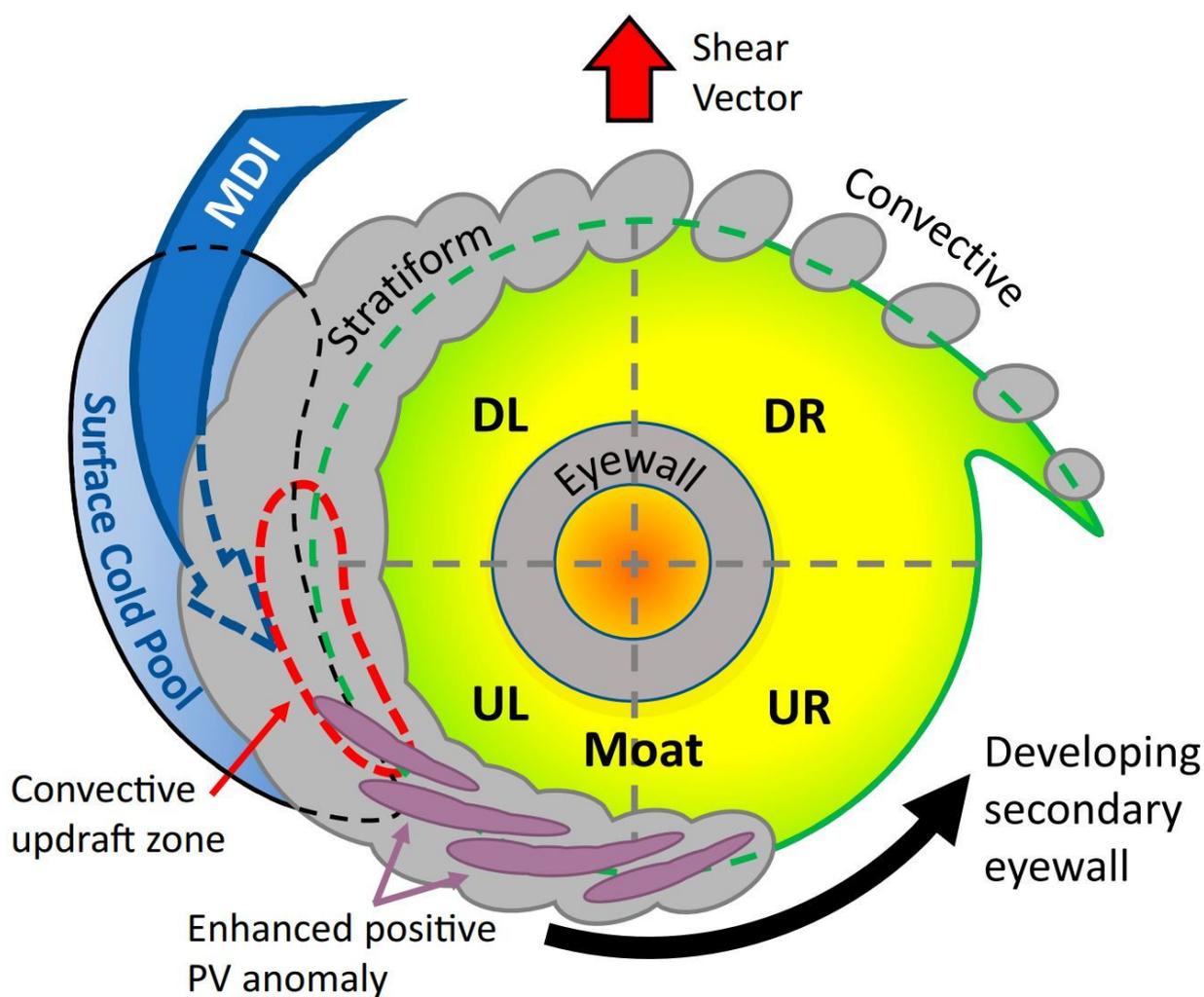
1206

1207 Fig. 8. Hurricane Matthew's track (solid line) and reflectivity at $z = 4$ km at 1930 UTC 6 Oct

1208 and 2126 7 Oct, respectively. Figure 1 of Cha et al. (2021). © American Meteorological

1209 Society. Used with permission.

1210

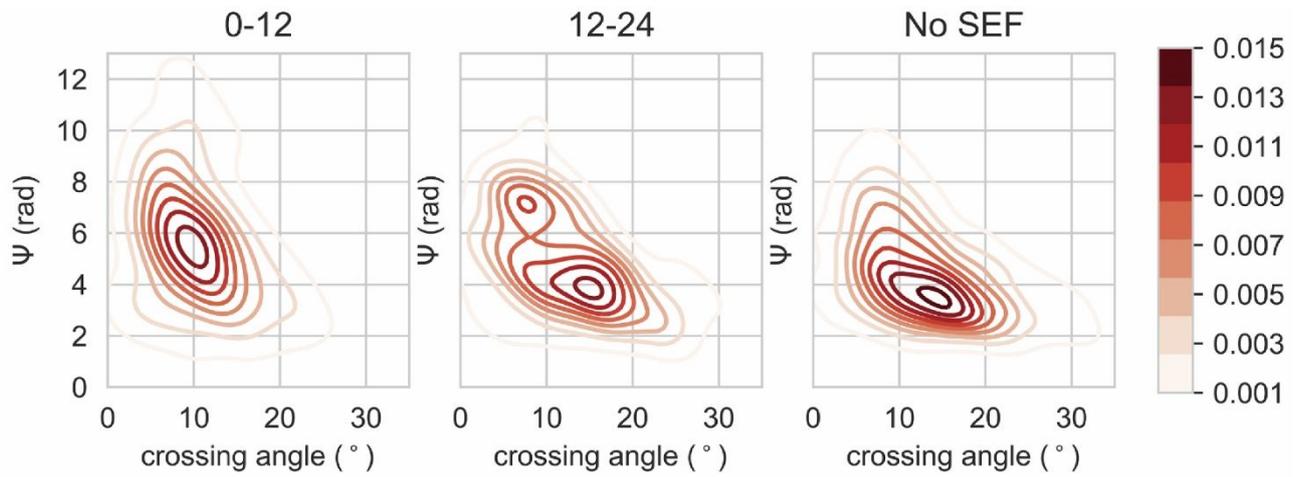


1211

1212 Fig. 9. A schematic diagram showing the SBC reflectivity (gray and purple) during the SEF
 1213 process. The MDI emerges at the left-of-shear quadrants, where the surface cold pool
 1214 underneath interacts with the high- θ_e envelope (yellow). Convective updrafts are
 1215 reinvigorated in this region, which then generates enhanced positive PV anomalies that
 1216 propagate cyclonically downwind to form the secondary eyewall. Figure 13 of Yu et al. (2022).

1217 © American Meteorological Society. Used with permission.

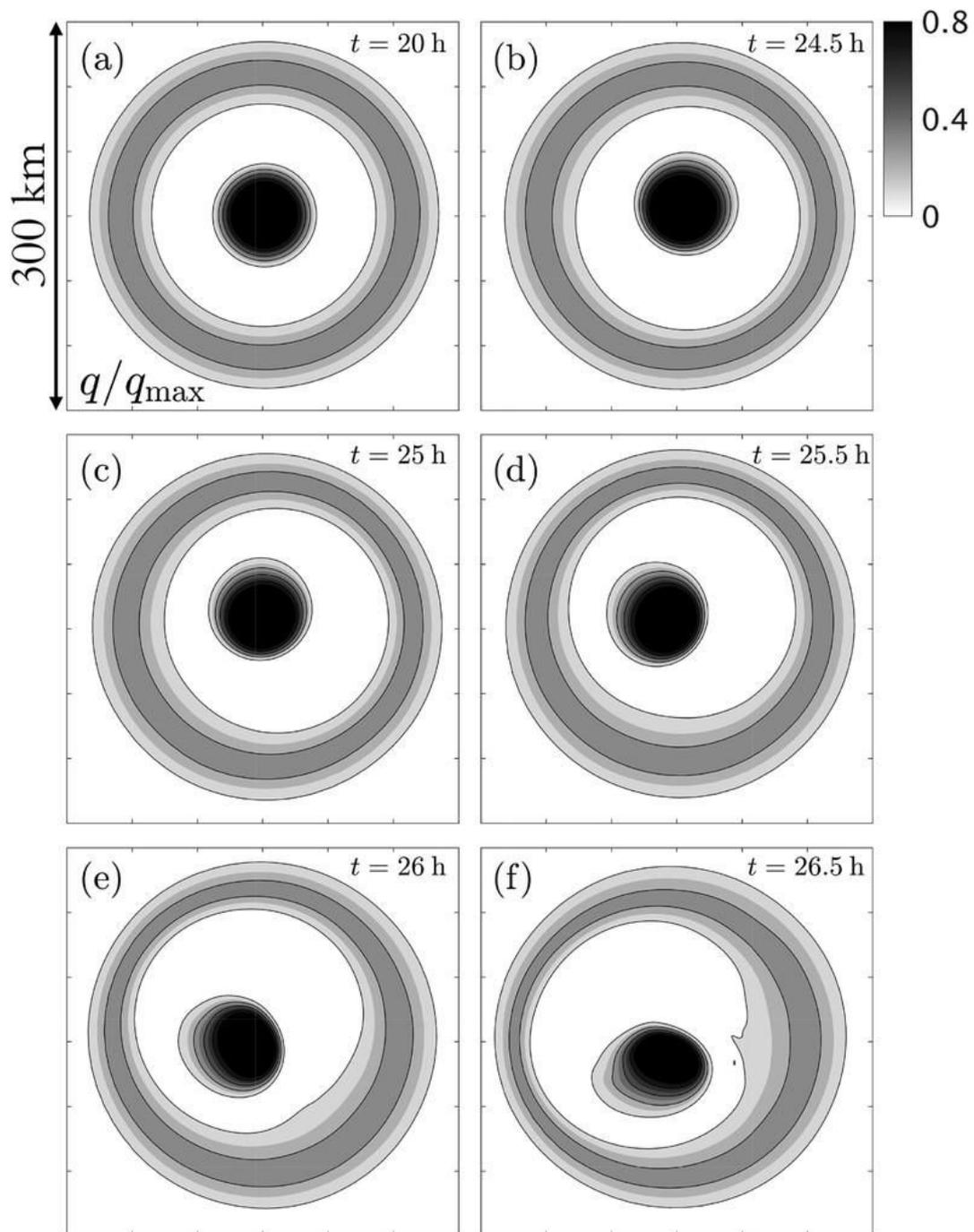
1218



1220

1221 Fig. 10. Kernel density estimation of SBC's azimuthal extent (Ψ) and crossing angle (a
1222 smaller value indicates higher circularity) prior to SEF (a-b), and (c) cases with no SEF within
1223 24 h. Figure 9 of Vaughan et al. (2020).

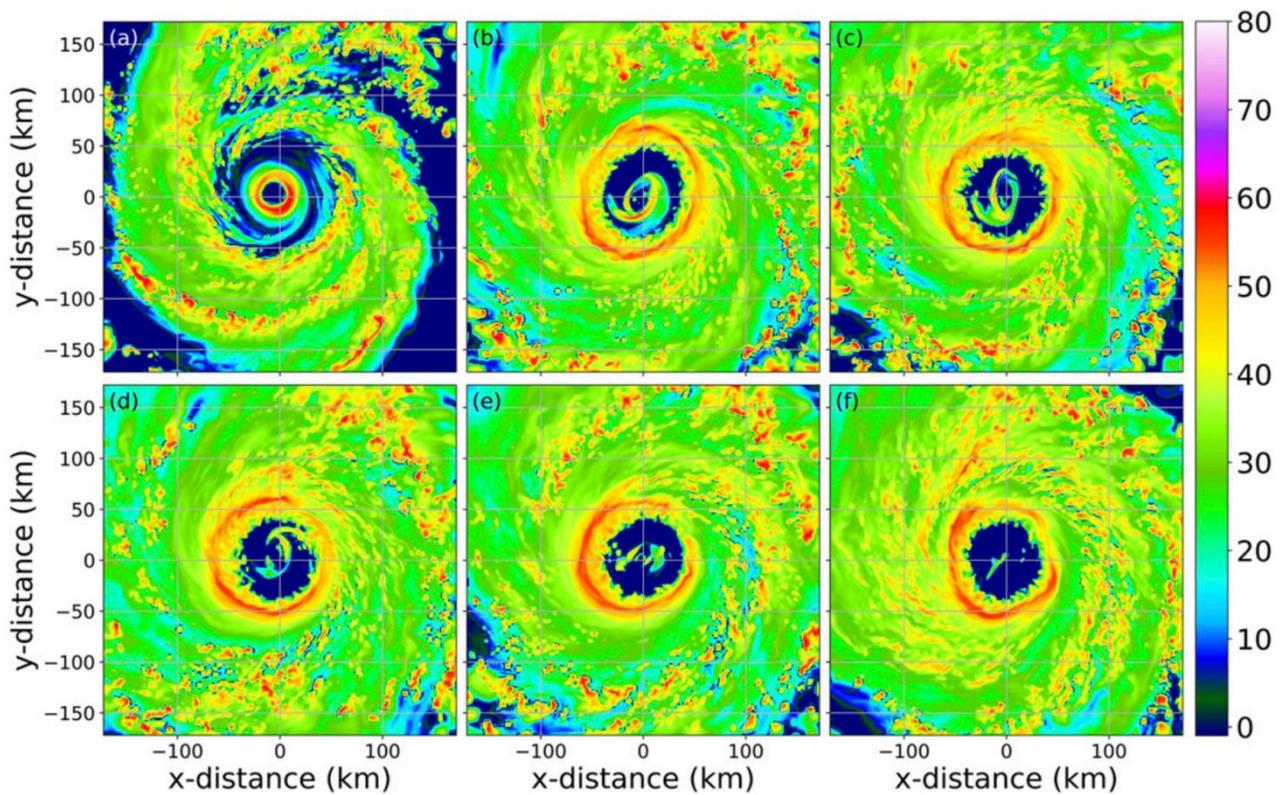
1224



1225

1226 Fig. 11. Normalized potential vorticity representing multiple eyewalls. Figure 3 of Menelaou

1227 et al. (2018). © American Meteorological Society. Used with permission.



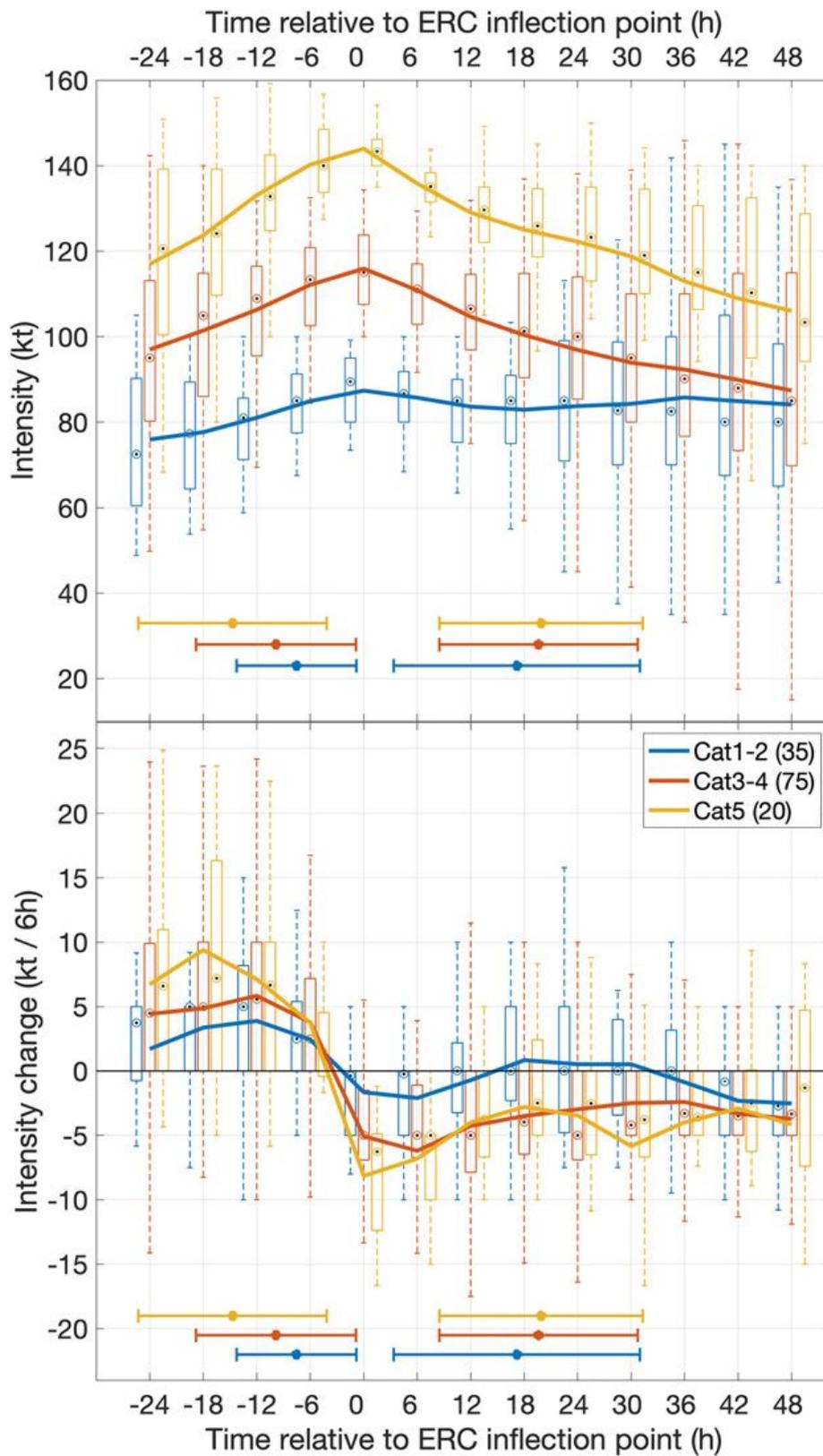
1228

1229 Fig. 12. The simulated radar reflectivity (dBZ) of Hurricane Wilma at $z = 2$ km at (a) $t = 42$,

1230 (b) 59.5, (c) 62, (d) 64, (e) 66, and (f) 69 h, exhibiting the elliptical inner eyewall structure.

1231 Figure 2 of Lai et al. (2019). © American Meteorological Society. Used with permission.

1232

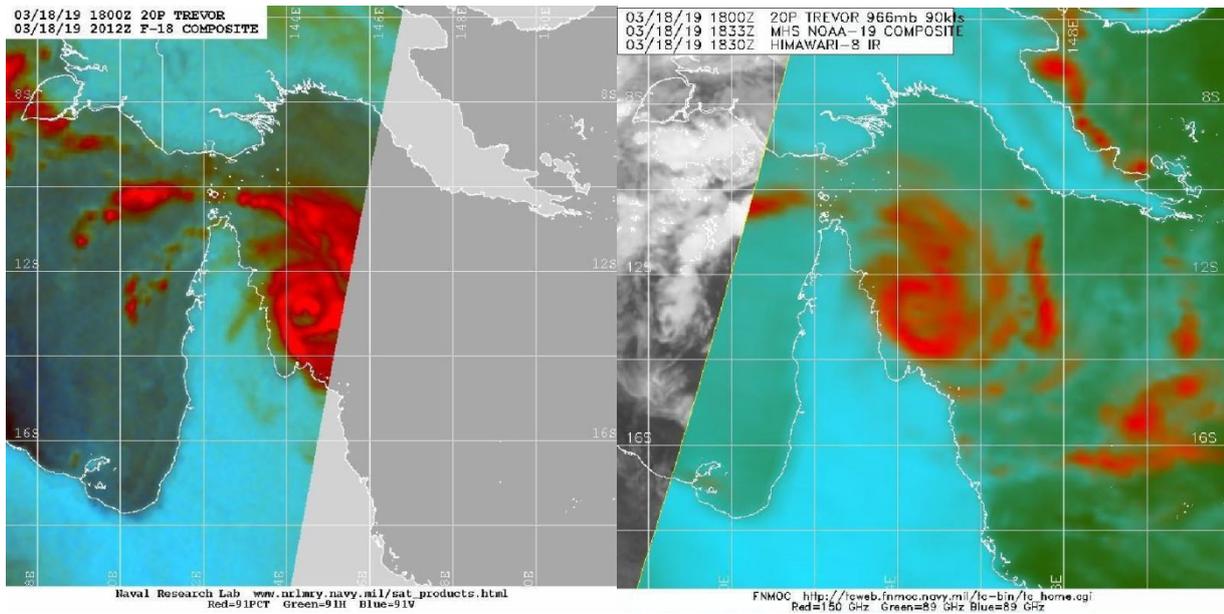


1233

1234 Fig. 13. (top) Intensity and (bottom) intensity change during Atlantic hurricane ERC events.

1235 See Kossin et al. (2023) for details. Fig. 3 of Kossin et al. (2023). © American Meteorological

1236 Society. Used with permission.

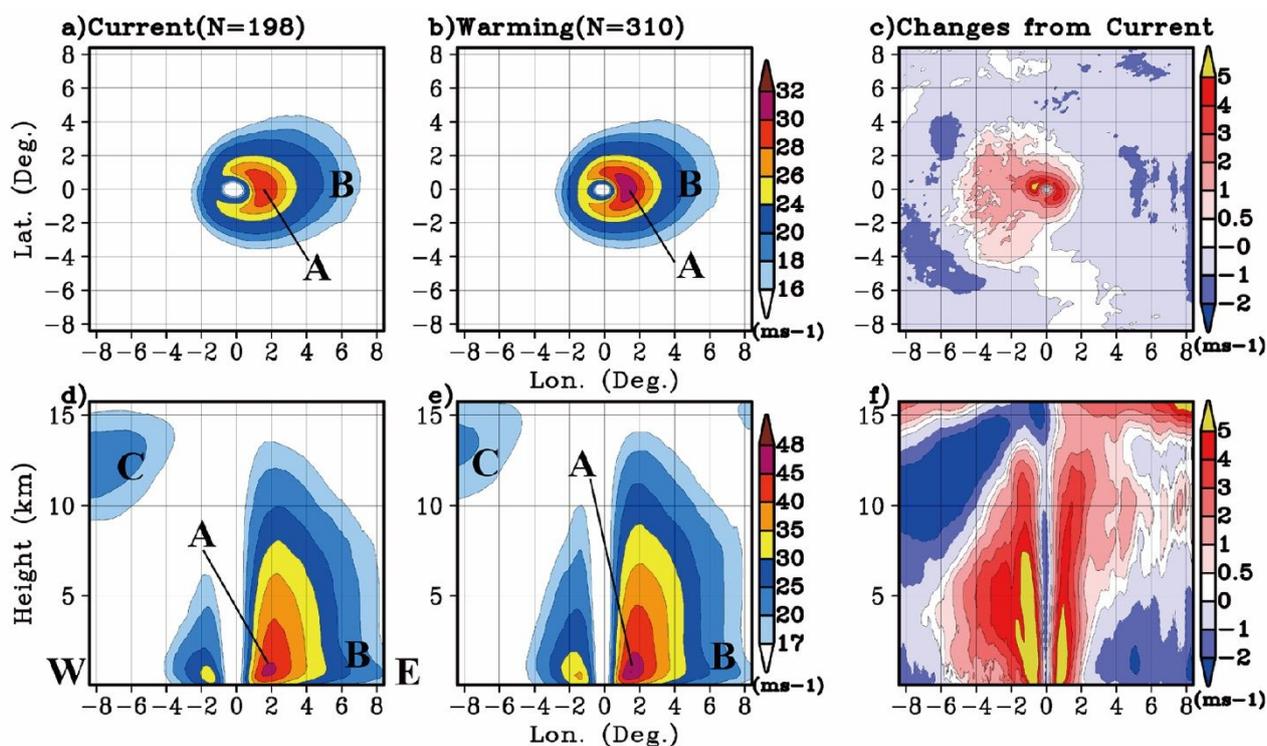


1237

1238 Fig. 14. Microwave images of Tropical Cyclone Trevor (2019) revealing double eyewall

1239 structure. Images courtesy of NRL.

1240



1241

1242 Fig. 15. Composite of TC-centered 10-m wind fields when the TC center was located in
 1243 142°E – 147°E and 35°N – 40°N in the (a) current and (b) warming climates. The differences
 1244 with respect to the current climate is shown in (c). (d)–(f) Same as (a)–(c), but for vertical
 1245 cross sections of horizontal wind speeds. ‘A’, ‘B’, and ‘C’ in the panels indicate the areas of
 1246 high winds, moderate wind speeds, and the jet, respectively. Figure 3 of Kanada et al. (2020).

1247

1248 **References**

- 1249 Abarca, S. F., and M. T. Montgomery, 2013: Essential dynamics of secondary eyewall
1250 formation. *J. Atmos. Sci.*, **70**, 3216-3230.
- 1251 Abarca, S. F., M. T. Montgomery, S. A. Braun, and J. Dunion, 2016: On the secondary
1252 eyewall formation of Hurricane Edouard (2014). *Mon. Wea. Rev.*, **144**, 3321-3331.
- 1253 Ahern, K., R. E. Hart, and M. A. Bourassa, 2021: Asymmetric hurricane boundary layer
1254 structure during storm decay. Part I: Formation of descending inflow. *Mon. Wea. Rev.*, **149**,
1255 3851-3874.
- 1256 —, 2022: Asymmetric Hurricane Boundary Layer Structure during Storm Decay. Part II:
1257 Secondary Eyewall Formation. *Mon. Wea. Rev.*, **150**, 1915-1936.
- 1258 Ahern, K., M. A. Bourassa, R. E. Hart, J. A. Zhang, and R. F. Rogers, 2019: Observed
1259 kinematic and thermodynamic structure in the hurricane boundary layer during intensity
1260 change. *Mon. Wea. Rev.*, **147**, 2765-2785.
- 1261 Aijaz, S., J. D. Kepert, H. Ye, Z. Huang, and A. Hawksford, 2019: Bias correction of tropical
1262 cyclone parameters in the ECMWF ensemble prediction system in Australia. *Mon. Wea.*
1263 *Rev.*, **147**, 4261-4285.
- 1264 Alvey, G. R., E. Zipser, and J. Zawislak, 2020: How does Hurricane Edouard (2014) evolve
1265 toward symmetry before rapid intensification? A high-resolution ensemble study. *J. Atmos.*
1266 *Sci.*, **77**, 1329-1351.
- 1267 Alvey III, G. R., and A. Hazelton, 2022: How Do Weak, Misaligned Tropical Cyclones Evolve
1268 Toward Alignment? A Multi - Case Study Using the Hurricane Analysis and Forecast
1269 System. *J. Geophys. Res. Atmos.*, **127**, e2022JD037268.
- 1270 Avenas, A., and Coauthors, 2023: Reexamining the estimation of tropical cyclone radius of
1271 maximum wind from outer size with an extensive synthetic aperture radar dataset. *Mon.*
1272 *Wea. Rev.*, **151**, 3169-3189.
- 1273 Carrasco, C. A., C. W. Landsea, and Y.-L. Lin, 2014: The influence of tropical cyclone size
1274 on its intensification. *Wea. Fore.*, **29**, 582-590.
- 1275 Cha, T.-Y., M. M. Bell, and A. J. DesRosiers, 2021: Doppler Radar Analysis of the Eyewall
1276 Replacement Cycle of Hurricane Matthew (2016) in Vertical Wind Shear. *Mon. Wea. Rev.*,
1277 **149**, 2927-2943.
- 1278 Chavas, D. R., and N. Lin, 2016: A model for the complete radial structure of the tropical
1279 cyclone wind field. Part II: Wind field variability. *J. Atmos. Sci.*, **73**, 3093-3113.
- 1280 Chavas, D. R., and J. A. Knaff, 2022: A simple model for predicting the tropical cyclone
1281 radius of maximum wind from outer size. *Wea. Fore.*, **37**, 563-579.
- 1282 Chavas, D. R., N. Lin, and K. Emanuel, 2015: A model for the complete radial structure of
1283 the tropical cyclone wind field. Part I: Comparison with observed structure. *J. Atmos. Sci.*,
1284 **72**, 3647-3662.
- 1285 Chen, G., 2018: Secondary eyewall formation and concentric eyewall replacement in
1286 association with increased low-level inner-core diabatic cooling. *J. Atmos. Sci.*, **75**, 2659-
1287 2685.
- 1288 Chen, J., and D. R. Chavas, 2020: The transient responses of an axisymmetric tropical
1289 cyclone to instantaneous surface roughening and drying. *J. Atmos. Sci.*, **77**, 2807-2834.
- 1290 Chen, X., M. Xue, and J. Fang, 2018: Rapid intensification of Typhoon Mujigae (2015) under
1291 different sea surface temperatures: Structural changes leading to rapid intensification. *J.*
1292 *Atmos. Sci.*, **75**, 4313-4335.
- 1293 Chen, X., J.-F. Gu, J. A. Zhang, F. D. Marks, R. F. Rogers, and J. J. Cione, 2021: Boundary
1294 layer recovery and precipitation symmetrization preceding rapid intensification of tropical
1295 cyclones under shear. *J. Atmos. Sci.*, **78**, 1523-1544.
- 1296 Chen, X., and Coauthors, 2023: Research advances on internal processes affecting tropical
1297 cyclone intensity change from 2018–2022. *Tropical Cyclone Research and Review*, **12**, 10-

- 1298 29.
- 1299 Cheng, C.-J., and C.-C. Wu, 2018: The role of WISHE in secondary eyewall formation. *J.*
- 1300 *Atmos. Sci.*, **75**, 3823-3841.
- 1301 Cheung, A. A., C. J. Slocum, J. A. Knaff, and M. N. Razin, 2024: Documenting the
- 1302 Progressions of Secondary Eyewall Formations. *Wea. Fore.*, **39**, 19-40.
- 1303 Corbosiero, K. L., and J. Molinari, 2003: The relationship between storm motion, vertical
- 1304 wind shear, and convective asymmetries in tropical cyclones. *J. Atmos. Sci.*, **60**, 366-376.
- 1305 Cronin, T. W., and D. R. Chavas, 2019: Dry and semidry tropical cyclones. *J. Atmos. Sci.*,
- 1306 **76**, 2193-2212.
- 1307 Dai, H., and Coauthors, 2021: Quasi - Periodic Intensification of Convective Asymmetries
- 1308 in the Outer Eyewall of Typhoon Lekima (2019). *Geophys. Res. Lett.*, **48**, e2020GL091633.
- 1309 DeMaria, M., M. Mainelli, L. K. Shay, J. A. Knaff, and J. Kaplan, 2005: Further improvements
- 1310 to the statistical hurricane intensity prediction scheme (SHIPS). *Wea. Fore.*, **20**, 531-543.
- 1311 DesRosiers, A. J., M. M. Bell, and T.-Y. Cha, 2022: Vertical vortex development in Hurricane
- 1312 Michael (2018) during rapid intensification. *Mon. Wea. Rev.*, **150**, 99-114.
- 1313 Didlake, A. C., and R. A. Houze, 2011: Kinematics of the secondary eyewall observed in
- 1314 Hurricane Rita (2005). *J. Atmos. Sci.*, **68**, 1620-1636.
- 1315 Didlake, A. C., G. M. Heymsfield, P. D. Reasor, and S. R. Guimond, 2017: Concentric
- 1316 eyewall asymmetries in Hurricane Gonzalo (2014) observed by airborne radar. *Mon. Wea.*
- 1317 *Rev.*, **145**, 729-749.
- 1318 Didlake, A. C., P. D. Reasor, R. F. Rogers, and W.-C. Lee, 2018: Dynamics of the transition
- 1319 from spiral rainbands to a secondary eyewall in Hurricane Earl (2010). *J. Atmos. Sci.*, **75**,
- 1320 2909-2929.
- 1321 Dvorak, V. F., 1984: *Tropical cyclone intensity analysis using satellite data*. Vol. 11, US
- 1322 Department of Commerce, National Oceanic and Atmospheric Administration, National
- 1323 Environmental Satellite, Data, and Information Service.
- 1324 Emanuel, K., and R. Rotunno, 2011: Self-stratification of tropical cyclone outflow. Part I:
- 1325 Implications for storm structure. *J. Atmos. Sci.*, **68**, 2236-2249.
- 1326 Emanuel, K., C. DesAutels, C. Holloway, and R. Korty, 2004: Environmental control of
- 1327 tropical cyclone intensity. *J. Atmos. Sci.*, **61**, 843-858.
- 1328 Emanuel, K. A., 1983: On assessing local conditional symmetric instability from atmospheric
- 1329 soundings. *Mon. Wea. Rev.*, **111**, 2016-2033.
- 1330 —, 1986: An air-sea interaction theory for tropical cyclones. Part I: Steady-state
- 1331 maintenance. *J. Atmos. Sci.*, **43**, 585-605.
- 1332 Fei, R., Y. Wang, and Y. Li, 2021: Contribution of vertical advection to supergradient wind
- 1333 in tropical cyclone boundary layer: A numerical study. *J. Atmos. Sci.*, **78**, 1057-1073.
- 1334 Fischer, M. S., R. F. Rogers, and P. D. Reasor, 2020: The rapid intensification and eyewall
- 1335 replacement cycles of Hurricane Irma (2017). *Mon. Wea. Rev.*, **148**, 981-1004.
- 1336 Francis, A., and B. Strahl, 2021: Joint Typhoon warning center annual tropical cyclone report
- 1337 2020. <https://www.metoc.navy.mil/jtwc/products/atcr/2020atcr.pdf>. Last accessed, 8, 2022.
- 1338 Fujiwara, K., T. Takemi, and N. Mori, 2023: Response of Intensity and Structure of Typhoon
- 1339 Jebi (2018) before Landfall to 2-K and 4-K Warmed Future Climates in Dynamical
- 1340 Downscaling Experiments. *SOLA*.
- 1341 Ge, X., 2015: Impacts of environmental humidity on concentric eyewall structure.
- 1342 *Atmospheric Science Letters*, **16**, 273-278.
- 1343 Green, T., X. Wang, and X. Lu, 2022: Impact of assimilating ground-based and airborne
- 1344 radar observations for the analysis and prediction of the eyewall replacement cycle of
- 1345 hurricane matthew (2016) using the HWRF hybrid 3DENVAR system. *Mon. Wea. Rev.*, **150**,
- 1346 1157-1175.
- 1347 Guan, L., and X. Ge, 2018: How does tropical cyclone size affect the onset timing of
- 1348 secondary eyewall formation? *Journal of Meteorological Research*, **32**, 124-134.

- 1349 Guimond, S. R., P. D. Reasor, G. M. Heymsfield, and M. M. McLinden, 2020: The dynamics
1350 of vortex Rossby waves and secondary eyewall development in Hurricane Matthew (2016):
1351 New insights from radar measurements. *J. Atmos. Sci.*, **77**, 2349-2374.
- 1352 Gutmann, E. D., and Coauthors, 2018: Changes in hurricanes from a 13-yr convection-
1353 permitting pseudo-global warming simulation. *J. Climate*, **31**, 3643-3657.
- 1354 Hawkins, H. F., and D. T. Rubsam, 1968: Hurricane Hilda, 1964. *Monthly Weather Review*,
1355 **96**, 617-636.
- 1356 Hirano, S., K. Ito, H. Yamada, S. Tsujino, K. Tsuboki, and C.-C. Wu, 2022: Deep Eye Clouds
1357 in Tropical Cyclone Trami (2018) during T-PARCIID Dropsonde Observations. *J. Atmos. Sci.*,
1358 **79**, 683-703.
- 1359 Hlywiak, J., and D. S. Nolan, 2019: The influence of oceanic barrier layers on tropical
1360 cyclone intensity as determined through idealized, coupled numerical simulations. *J. Phys.*
1361 *Oceanogr.*, **49**, 1723-1745.
- 1362 —, 2021: The response of the near-surface tropical cyclone wind field to inland surface
1363 roughness length and soil moisture content during and after landfall. *J. Atmos. Sci.*, **78**, 983-
1364 1000.
- 1365 —, 2022: The evolution of asymmetries in the tropical cyclone boundary layer wind field
1366 during landfall. *Mon. Wea. Rev.*, **150**, 529-549.
- 1367 Horinouchi, T., S. Tsujino, M. Hayashi, U. Shimada, W. Yanase, A. Wada, and H. Yamada,
1368 2023: Stationary and transient asymmetric features in tropical cyclone eye with
1369 wavenumber-1 instability: Case study for Typhoon Haishen (2020) with atmospheric motion
1370 vectors from 30-second imaging. *Mon. Wea. Rev.*, **151**, 253-273.
- 1371 Huang, Y.-H., M. T. Montgomery, and C.-C. Wu, 2012: Concentric eyewall formation in
1372 Typhoon Sinlaku (2008). Part II: Axisymmetric dynamical processes. *J. Atmos. Sci.*, **69**, 662-
1373 674.
- 1374 Huang, Y.-H., C.-C. Wu, and M. T. Montgomery, 2018: Concentric eyewall formation in
1375 Typhoon Sinlaku (2008). Part III: Horizontal momentum budget analyses. *J. Atmos. Sci.*, **75**,
1376 3541-3563.
- 1377 Ito, J., T. Oizumi, and H. Niino, 2017: Near-surface coherent structures explored by large
1378 eddy simulation of entire tropical cyclones. *Scientific reports*, **7**, 3798.
- 1379 Ito, K., and R. Yamamoto, 2022: Thermodynamic and kinematic structure of tropical
1380 cyclones in the western North Pacific based on ACARS/AMDAR. *Frontiers in Earth Science*,
1381 **10**, 2354.
- 1382 Ito, K., Y. Ishikawa, Y. Miyamoto, and T. Awaji, 2011: Short-time-scale processes in a
1383 mature hurricane as a response to sea surface fluctuations. *J. Atmos. Sci.*, **68**, 2250-2272.
- 1384 John, E. B., K. Balaguru, L. R. Leung, G. R. Foltz, R. D. Hetland, and S. M. Hagos, 2023:
1385 Intensification of Hurricane Sally (2020) over the Mississippi River Plume. *Wea. Fore.*, **38**,
1386 1391-1404.
- 1387 Kanada, S., and A. Nishii, 2023: Observed concentric eyewalls of supertyphoon Hinnamnor
1388 (2022). *SOLA*, **19**, 70-77.
- 1389 Kanada, S., K. Tsuboki, and I. Takayabu, 2020: Future changes of tropical cyclones in the
1390 midlatitudes in 4-km-mesh downscaling experiments from large-ensemble simulations.
1391 *SOLA*, **16**, 57-63.
- 1392 Kasami, K., and M. Satoh, 2024: Mechanism of secondary eyewall formation in tropical
1393 cyclones revealed by sensitivity experiments on the mesoscale descending inflow. *SOLA*.
- 1394 Kepert, J., 2001: The dynamics of boundary layer jets within the tropical cyclone core. Part
1395 I: Linear theory. *J. Atmos. Sci.*, **58**, 2469-2484.
- 1396 Kepert, J., and Y. Wang, 2001: The dynamics of boundary layer jets within the tropical
1397 cyclone core. Part II: Nonlinear enhancement. *J. Atmos. Sci.*, **58**, 2485-2501.
- 1398 Kepert, J. D., 2013: How does the boundary layer contribute to eyewall replacement cycles
1399 in axisymmetric tropical cyclones? *J. Atmos. Sci.*, **70**, 2808-2830.

- 1400 —, 2017: Time and space scales in the tropical cyclone boundary layer, and the location of
1401 the eyewall updraft. *J. Atmos. Sci.*, **74**, 3305-3323.
- 1402 —, 2018: The boundary layer dynamics of tropical cyclone rainbands. *J. Atmos. Sci.*, **75**,
1403 3777-3795.
- 1404 Kepert, J. D., and D. S. Nolan, 2014: Reply to “Comments on ‘How does the boundary layer
1405 contribute to eyewall replacement cycles in axisymmetric tropical cyclones?’”. *J. Atmos. Sci.*,
1406 **71**, 4692-4704.
- 1407 Knaff, J. A., C. R. Sampson, and B. R. Strahl, 2020: A tropical cyclone rapid intensification
1408 prediction aid for the Joint Typhoon Warning Center’s areas of responsibility. *Wea. Fore.*,
1409 **35**, 1173-1185.
- 1410 Knaff, J. A., and Coauthors, 2021: Estimating tropical cyclone surface winds: Current status,
1411 emerging technologies, historical evolution, and a look to the future. *Tropical Cyclone
1412 Research and Review*, **10**, 125-150.
- 1413 Knutson, T., and Coauthors, 2020: Tropical cyclones and climate change assessment: Part
1414 II: Projected response to anthropogenic warming. *Bull. Amer. Meteor. Soc.*, **101**, E303-E322.
- 1415 Kossin, J. P., and M. Sitkowski, 2009: An objective model for identifying secondary eyewall
1416 formation in hurricanes. *Mon. Wea. Rev.*, **137**, 876-892.
- 1417 Kossin, J. P., and M. DeMaria, 2016: Reducing operational hurricane intensity forecast
1418 errors during eyewall replacement cycles. *Wea. Fore.*, **31**, 601-608.
- 1419 Kossin, J. P., W. H. Schubert, and M. T. Montgomery, 2000: Unstable interactions between
1420 a hurricane’s primary eyewall and a secondary ring of enhanced vorticity. *J. Atmos. Sci.*, **57**,
1421 3893-3917.
- 1422 Kossin, J. P., D. C. Herndon, A. J. Wimmers, X. Guo, and E. S. Blake, 2023: M-PERC: A
1423 new satellite microwave-based model to diagnose the onset of tropical cyclone eyewall
1424 replacement cycles. *Wea. Fore.*, **38**, 1405-1411.
- 1425 Kossin, J. P., and Coauthors, 2007: Estimating hurricane wind structure in the absence of
1426 aircraft reconnaissance. *Wea. Forecasting*, **22**, 89-101.
- 1427 Kuo, H.-C., L. Lin, C. Chang, and R. Williams, 2004: The formation of concentric vorticity
1428 structures in typhoons. *J. Atmos. Sci.*, **61**, 2722-2734.
- 1429 Kuo, H.-C., C.-P. Chang, Y.-T. Yang, and H.-J. Jiang, 2009: Western North Pacific typhoons
1430 with concentric eyewalls. *Mon. Wea. Rev.*, **137**, 3758-3770.
- 1431 Kuo, H. C., S. Tsujino, T. Y. Hsu, M. S. Peng, and S. H. Su, 2022: Scaling law for boundary
1432 layer inner eyewall pumping in concentric eyewalls. *J. Geophys. Res. Atmos.*, **127**,
1433 e2021JD035518.
- 1434 Lai, T.-K., K. Menelaou, and M. K. Yau, 2019: Barotropic instability across the moat and
1435 inner eyewall dissipation: A numerical study of Hurricane Wilma (2005). *J. Atmos. Sci.*, **76**,
1436 989-1013.
- 1437 Lai, T.-K., E. A. Hendricks, and M. K. Yau, 2021a: Long-term effect of barotropic instability
1438 across the moat in double-eyewall tropical cyclone-like vortices in forced and unforced
1439 shallow-water models. *J. Atmos. Sci.*, **78**, 4103-4126.
- 1440 Lai, T.-K., E. A. Hendricks, K. Menelaou, and M. K. Yau, 2021b: Roles of barotropic
1441 instability across the moat in inner eyewall decay and outer eyewall intensification: Three-
1442 dimensional numerical experiments. *J. Atmos. Sci.*, **78**, 473-496.
- 1443 Lai, T.-K., E. A. Hendricks, M. K. Yau, and K. Menelaou, 2021c: Roles of barotropic instability
1444 across the moat in inner eyewall decay and outer eyewall intensification: Essential dynamics.
1445 *J. Atmos. Sci.*, **78**, 1411-1428.
- 1446 Lau, K.-H., C. Y. Tam, and C. C. Wu, 2024: Island - induced eyewall replacement in a
1447 landfalling tropical cyclone: A model study of Super Typhoon Mangkhut (2018). *J. Geophys.
1448 Res. Atmos.*, **129**, e2023JD039541.
- 1449 Le Hénaff, M., and Coauthors, 2021: The role of the Gulf of Mexico ocean conditions in the

- 1450 intensification of Hurricane Michael (2018). *J. Geophys. Res. Oceans*, **126**, e2020JC016969.
- 1451 Li, X., Z. Pu, and Z. Gao, 2021a: Effects of roll vortices on the evolution of Hurricane Harvey
- 1452 during landfall. *J. Atmos. Sci.*, **78**, 1847-1867.
- 1453 Li, X., X. Cheng, J. Fei, X. Huang, and J. Ding, 2022a: The modulation effect of sea surface
- 1454 cooling on the eyewall replacement cycle in Typhoon Trami (2018). *Mon. Wea. Rev.*, **150**,
- 1455 1417-1436.
- 1456 Li, Y., Y. Wang, and Y. Lin, 2019: Revisiting the dynamics of eyewall contraction of tropical
- 1457 cyclones. *J. Atmos. Sci.*, **76**, 3229-3245.
- 1458 Li, Y., Y. Wang, and Z.-M. Tan, 2022b: How frequently does rapid intensification occur after
- 1459 rapid contraction of the radius of maximum wind in tropical cyclones over the North Atlantic
- 1460 and Eastern North Pacific? *Mon. Wea. Rev.*, **150**, 1747-1760.
- 1461 Li, Y., Y. Wang, Y. Lin, and X. Wang, 2021b: Why does rapid contraction of the radius of
- 1462 maximum wind precede rapid intensification in tropical cyclones? *J. Atmos. Sci.*, **78**, 3441-
- 1463 3453.
- 1464 Liu, Q., L. Wu, N. Qin, and Y. Li, 2021: Storm - scale and fine - scale boundary layer
- 1465 structures of tropical cyclones simulated with the WRF - LES framework. *J. Geophys. Res.*
- 1466 *Atmos.*, **126**, e2021JD035511.
- 1467 Liu, X., Q. Li, and Y. Dai, 2022: Stronger vertical shear leads to earlier secondary eyewall
- 1468 formation in idealized numerical simulations. *Geophys. Res. Lett.*, **49**, e2022GL098093.
- 1469 Macdonald, N. J., 1968: The evidence for the existence of Rossby - like waves in the
- 1470 hurricane vortex. *Tellus*, **20**, 138-150.
- 1471 Mallen, K. J., M. T. Montgomery, and B. Wang, 2005: Reexamining the near-core radial
- 1472 structure of the tropical cyclone primary circulation: Implications for vortex resiliency. *J.*
- 1473 *Atmos. Sci.*, **62**, 408-425.
- 1474 Martinez, J., C. C. Nam, and M. M. Bell, 2020: On the contributions of incipient vortex
- 1475 circulation and environmental moisture to tropical cyclone expansion. *J. Geophys. Res.*
- 1476 *Atmos.*, **125**, e2020JD033324.
- 1477 Menelaou, K., M. K. Yau, and Y. Martinez, 2012: On the dynamics of the secondary eyewall
- 1478 genesis in Hurricane Wilma (2005). *Geophys. Res. Lett.*, **39**.
- 1479 Menelaou, K., M. K. Yau, and T.-K. Lai, 2018: A Possible Three-Dimensional Mechanism
- 1480 for Oscillating Wobbles in Tropical Cyclone-Like Vortices with Concentric Eyewalls. *J.*
- 1481 *Atmos. Sci.*, **75**, 2157-2174.
- 1482 Miura, H., and Coauthors, 2023: Asymptotic matching between weather and climate models.
- 1483 *Bull. Amer. Meteor. Soc.*, **104**, E2308-E2315.
- 1484 Miyamoto, Y., D. S. Nolan, and N. Sugimoto, 2018: A dynamical mechanism for secondary
- 1485 eyewall formation in tropical cyclones. *J. Atmos. Sci.*, **75**, 3965-3986.
- 1486 Miyamoto, Y., H. Fudeyasu, and A. Wada, 2022: Intensity and structural changes of
- 1487 numerically simulated Typhoon Faxai (1915) before landfall. *Journal of the Meteorological*
- 1488 *Society of Japan. Ser. II*, **100**, 181-196.
- 1489 Molinari, J., J. A. Zhang, R. F. Rogers, and D. Vollaro, 2019: Repeated eyewall replacement
- 1490 cycles in Hurricane Frances (2004). *Mon. Wea. Rev.*, **147**, 2009-2022.
- 1491 Montgomery, M. T., and R. J. Kallenbach, 1997: A theory for vortex Rossby-waves and its
- 1492 application to spiral bands and intensity changes in hurricanes. *Quart. J. Roy. Meteor. Soc.*,
- 1493 **123**, 435-466.
- 1494 Moon, Y., and D. S. Nolan, 2010: The dynamic response of the hurricane wind field to spiral
- 1495 rainband heating. *J. Atmos. Sci.*, **67**, 1779-1805.
- 1496 Moon, Y., D. S. Nolan, and M. Iskandarani, 2010: On the use of two-dimensional
- 1497 incompressible flow to study secondary eyewall formation in tropical cyclones. *J. Atmos.*
- 1498 *Sci.*, **67**, 3765-3773.
- 1499 Mrowiec, A. A., S. T. Garner, and O. M. Pauluis, 2011: Axisymmetric hurricane in a dry

- 1500 atmosphere: Theoretical framework and numerical experiments. *J. Atmos. Sci.*, **68**, 1607-
1501 1619.
- 1502 Nolan, D. S., M. T. Montgomery, and L. D. Grasso, 2001: The wavenumber-one instability
1503 and trochoidal motion of hurricane-like vortices. *J. Atmos. Sci.*, **58**, 3243-3270.
- 1504 Nong, S., and K. Emanuel, 2003: A numerical study of the genesis of concentric eyewalls in
1505 hurricanes. *Quarterly Journal of the Royal Meteorological Society: A journal of the
1506 atmospheric sciences, applied meteorology and physical oceanography*, **129**, 3323-3338.
- 1507 Ooyama, K., 1969: Numerical simulation of the life cycle of tropical cyclones. *J. Atmos. Sci.*,
1508 **26**, 3-40.
- 1509 Patricola, C. M., and M. F. Wehner, 2018: Anthropogenic influences on major tropical
1510 cyclone events. *Nature*, **563**, 339-346.
- 1511 Peng, K., R. Rotunno, G. H. Bryan, and J. Fang, 2019: Evolution of an axisymmetric tropical
1512 cyclone before reaching slantwise moist neutrality. *J. Atmos. Sci.*, **76**, 1865-1884.
- 1513 Persing, J., and M. T. Montgomery, 2022: Does the Rotating Convection Paradigm Describe
1514 Secondary Eyewall Formation in Idealized Three-Dimensional Simulations? *J. Atmos. Sci.*,
1515 **79**, 745-762.
- 1516 Powell, M. D., P. J. Vickery, and T. A. Reinhold, 2003: Reduced drag coefficient for high
1517 wind speeds in tropical cyclones. *Nature*, **422**, 279-283.
- 1518 Pulmano, L., and L. Joykuty, 2021: Comparison of ARCHER MPERC to NHC Analysis.
1519 *Journal of Student Research*, **10**.
- 1520 Qin, N., and L. Wu, 2021: Possible environmental influence on eyewall expansion during the
1521 rapid intensification of hurricane Helene (2006). *Frontiers in Earth Science*, **9**, 715012.
- 1522 Qin, N., L. Wu, and Q. Liu, 2021: Evolution of the moat associated with the secondary
1523 eyewall formation in a simulated tropical cyclone. *J. Atmos. Sci.*, **78**, 4021-4035.
- 1524 Qiu, X., Z.-M. Tan, and Q. Xiao, 2010: The roles of vortex Rossby waves in hurricane
1525 secondary eyewall formation. *Mon. Wea. Rev.*, **138**, 2092-2109.
- 1526 Razin, M. N., and M. M. Bell, 2021: The unconventional eyewall replacement cycle of
1527 Hurricane Ophelia (2005). *Mon. Wea. Rev.*, **149**, 2151-2170.
- 1528 Reasor, P. D., M. T. Montgomery, and L. D. Grasso, 2004: A new look at the problem of
1529 tropical cyclones in vertical shear flow: Vortex resiliency. *J. Atmos. Sci.*, **61**, 3-22.
- 1530 Rios-Berrios, R., C. A. Davis, and R. D. Torn, 2018: A hypothesis for the intensification of
1531 tropical cyclones under moderate vertical wind shear. *J. Atmos. Sci.*, **75**, 4149-4173.
- 1532 Rogers, R., P. Reasor, and S. Lorsolo, 2013: Airborne Doppler observations of the inner-
1533 core structural differences between intensifying and steady-state tropical cyclones. *Mon.
1534 Wea. Rev.*, **141**, 2970-2991.
- 1535 Rogers, R. F., and J. A. Zhang, 2023: Airborne Doppler radar observations of tropical
1536 cyclone boundary layer kinematic structure and evolution during landfall. *Geophys. Res.
1537 Lett.*, **50**, e2023GL105548.
- 1538 Rogers, R. F., and Coauthors, 2017: Rewriting the tropical record books: The extraordinary
1539 intensification of Hurricane Patricia (2015). *Bull. Amer. Meteor. Soc.*, **98**, 2091-2112.
- 1540 Rostami, M., and V. Zeitlin, 2022: Evolution of double-eye wall hurricanes and emergence
1541 of complex tripolar end states in moist-convective rotating shallow water model. *Physics of
1542 Fluids*, **34**.
- 1543 Rozoff, C. M., W. H. Schubert, B. D. McNoldy, and J. P. Kossin, 2006: Rapid filamentation
1544 zones in intense tropical cyclones. *J. Atmos. Sci.*, **63**, 325-340.
- 1545 Rudzin, J. E., L. K. Shay, and W. E. Johns, 2018: The influence of the barrier layer on SST
1546 response during tropical cyclone wind forcing using idealized experiments. *J. Phys.
1547 Oceanogr.*, **48**, 1471-1478.
- 1548 Shapiro, L. J., 1983: The asymmetric boundary layer flow under a translating hurricane.
1549 *Journal of Atmospheric Sciences*, **40**, 1984-1998.
- 1550 Shen, L.-Z., C.-C. Wu, and F. Judt, 2021: The role of surface heat fluxes on the size of

- 1551 Typhoon Megi (2016). *J. Atmos. Sci.*, **78**, 1075-1093.
- 1552 Simpson, R. H., 1952: Exploring eye of Typhoon "Marge," 1951. *Bull. Amer. Meteor. Soc.*,
- 1553 **33**, 286-298.
- 1554 Sitkowski, M., J. P. Kossin, and C. M. Rozoff, 2011: Intensity and structure changes during
- 1555 hurricane eyewall replacement cycles. *Mon. Wea. Rev.*, **139**, 3829-3847.
- 1556 Slocum, C. J., J. A. Knaff, and S. N. Stevenson, 2023a: Lightning-based tropical cyclone
- 1557 rapid intensification guidance. *Wea. Fore.*, **38**, 1209-1227.
- 1558 Slocum, C. J., R. K. Taft, J. P. Kossin, and W. H. Schubert, 2023b: Barotropic instability
- 1559 during eyewall replacement. *Meteorology*, **2**, 191-221.
- 1560 Song, C., S. Park, and J. Shin, 2020: Tropical Cyclone Activities in Warm Climate with
- 1561 Quadrupled CO₂ Concentration Simulated by a New General Circulation Model. *J. Geophys.*
- 1562 *Res. Atmos.*, **125**, e2019JD032314.
- 1563 Stern, D., and Coauthors, 2014a: Subtopic 2.6. Intensity change: Internal processes. Report
- 1564 for the Eighth WMO International Workshop on Tropical Cyclones IWTC-VIII. *World*
- 1565 *Meteorological Organization, Jeju Island, South Korea.*
- 1566 Stern, D. P., and F. Zhang, 2013: How does the eye warm? Part I: A potential temperature
- 1567 budget analysis of an idealized tropical cyclone. *J. Atmos. Sci.*, **70**, 73-90.
- 1568 Stern, D. P., J. R. Brisbois, and D. S. Nolan, 2014b: An expanded dataset of hurricane
- 1569 eyewall sizes and slopes. *J. Atmos. Sci.*, **71**, 2747-2762.
- 1570 Stern, D. P., J. L. Vigh, D. S. Nolan, and F. Zhang, 2015: Revisiting the relationship between
- 1571 eyewall contraction and intensification. *J. Atmos. Sci.*, **72**, 1283-1306.
- 1572 Stern, D. P., J. D. Kepert, G. H. Bryan, and J. D. Doyle, 2020: Understanding atypical
- 1573 midlevel wind speed maxima in hurricane eyewalls. *J. Atmos. Sci.*, **77**, 1531-1557.
- 1574 Tao, D., and F. Zhang, 2019: Evolution of dynamic and thermodynamic structures before
- 1575 and during rapid intensification of tropical cyclones: Sensitivity to vertical wind shear. *Mon.*
- 1576 *Wea. Rev.*, **147**, 1171-1191.
- 1577 Tao, D., R. Rotunno, and M. Bell, 2020: Lilly's model for steady-state tropical cyclone
- 1578 intensity and structure. *J. Atmos. Sci.*, **77**, 3701-3720.
- 1579 Terwey, W. D., and M. T. Montgomery, 2008: Secondary eyewall formation in two idealized,
- 1580 full - physics modeled hurricanes. *J. Geophys. Res. Atmos.*, **113**.
- 1581 Trabling, B. C., and M. M. Bell, 2021: The sensitivity of eyewall replacement cycles to
- 1582 shortwave radiation. *J. Geophys. Res. Atmos.*, **126**, e2020JD034016.
- 1583 Tsujino, S., T. Horinouchi, T. Tsukada, H. C. Kuo, H. Yamada, and K. Tsuboki, 2021: Inner
- 1584 - core wind field in a concentric eyewall replacement of Typhoon Trami (2018): A
- 1585 quantitative analysis based on the Himawari - 8 satellite. *J. Geophys. Res. Atmos.*, **126**,
- 1586 e2020JD034434.
- 1587 Tsukada, T., and T. Horinouchi, 2020: Estimation of the tangential winds and asymmetric
- 1588 structures in typhoon inner core region using Himawari - 8. *Geophys. Res. Lett.*, **47**,
- 1589 e2020GL087637.
- 1590 —, 2023: Strong Relationship between Eye Radius and Radius of Maximum Wind of
- 1591 Tropical Cyclones. *Mon. Wea. Rev.*, **151**, 569-588.
- 1592 Ueno, M., and M. Kunii, 2009: Some aspects of azimuthal wavenumber-one structure of
- 1593 typhoons represented in the JMA operational mesoscale analyses. *Journal of the*
- 1594 *Meteorological Society of Japan. Ser. II*, **87**, 615-633.
- 1595 Uhlhorn, E. W., B. W. Klotz, T. Vukicevic, P. D. Reasor, and R. F. Rogers, 2014: Observed
- 1596 hurricane wind speed asymmetries and relationships to motion and environmental shear.
- 1597 *Mon. Wea. Rev.*, **142**, 1290-1311.
- 1598 Vaughan, A., K. J. Walsh, and J. D. Kepert, 2020: The stationary banding complex and
- 1599 secondary eyewall formation in tropical cyclones. *J. Geophys. Res. Atmos.*, **125**,
- 1600 e2019JD031515.

- 1601 Vigh, J. L., and Coauthors, 2018: Topic 3.1. Intensity change: internal influences. Report of
1602 the Ninth WMO International Workshop on Tropical Cyclones (IWTC-9). *World*
1603 *Meteorological Organization, Honolulu, Hawaii*.
- 1604 Wadler, J. B., R. F. Rogers, and P. D. Reasor, 2018: The relationship between spatial
1605 variations in the structure of convective bursts and tropical cyclone intensification as
1606 determined by airborne Doppler radar. *Mon. Wea. Rev.*, **146**, 761-780.
- 1607 Wadler, J. B., D. S. Nolan, J. A. Zhang, and L. K. Shay, 2021a: Thermodynamic
1608 characteristics of downdrafts in tropical cyclones as seen in idealized simulations of different
1609 intensities. *J. Atmos. Sci.*, **78**, 3503-3524.
- 1610 Wadler, J. B., J. A. Zhang, R. F. Rogers, B. Jaimes, and L. K. Shay, 2021b: The rapid
1611 intensification of Hurricane Michael (2018): Storm structure and the relationship to
1612 environmental and air-sea interactions. *Mon. Wea. Rev.*, **149**, 245-267.
- 1613 Wang, D., and Y. Lin, 2020: Size and structure of dry and moist reversible tropical cyclones.
1614 *J. Atmos. Sci.*, **77**, 2091-2114.
- 1615 Wang, H., and Y. Wang, 2021: A numerical study of Typhoon Megi (2010). Part II: Eyewall
1616 evolution crossing the Luzon Island. *Mon. Wea. Rev.*, **149**, 375-394.
- 1617 Wang, H., Y. Wang, J. Xu, and Y. Duan, 2019a: The axisymmetric and asymmetric aspects
1618 of the secondary eyewall formation in a numerically simulated tropical cyclone under
1619 idealized conditions on an f plane. *J. Atmos. Sci.*, **76**, 357-378.
- 1620 —, 2019b: Evolution of the warm-core structure during the eyewall replacement cycle in a
1621 numerically simulated tropical cyclone. *J. Atmos. Sci.*, **76**, 2559-2573.
- 1622 Wang, X., Q. Li, and N. E. Davidson, 2018: The coupled dynamic and thermodynamic
1623 processes for secondary eyewall formation. *J. Geophys. Res. Atmos.*, **123**, 9192-9219.
- 1624 Wang, Y.-F., and Z.-M. Tan, 2020: Outer rainbands-driven secondary eyewall formation of
1625 tropical cyclones. *J. Atmos. Sci.*, **77**, 2217-2236.
- 1626 —, 2022: Essential dynamics of the vertical wind shear affecting the secondary eyewall
1627 formation in tropical cyclones. *J. Atmos. Sci.*, **79**, 2831-2847.
- 1628 Wehner, M., 2021: Simulated changes in tropical cyclone size, accumulated cyclone energy
1629 and power dissipation index in a warmer climate. *Oceans*, MDPI, 688-699.
- 1630 Willoughby, H., J. Clos, and M. Shoreibah, 1982: Concentric eye walls, secondary wind
1631 maxima, and the evolution of the hurricane vortex. *J. Atmos. Sci.*, **39**, 395-411.
- 1632 Willoughby, H. E., F. D. Marks, and R. J. Feinberg, 1984: Stationary and moving convective
1633 bands in hurricanes. *J. Atmos. Sci.*, **41**, 3189-3211.
- 1634 Wood, K., and Coauthors, 2023: Phase transitions between tropical, subtropical, and
1635 extratropical cyclones: A review from IWTC-10. *Tropical Cyclone Research and Review*.
- 1636 Wu, C.-C., Y.-H. Huang, and G.-Y. Lien, 2012: Concentric eyewall formation in Typhoon
1637 Sinlaku (2008). Part I: Assimilation of T-PARC data based on the ensemble Kalman filter
1638 (EnKF). *Mon. Wea. Rev.*, **140**, 506-527.
- 1639 Wu, Q., and Z. Ruan, 2021: Rapid contraction of the radius of maximum tangential wind and
1640 rapid intensification of a tropical cyclone. *J. Geophys. Res. Atmos.*, **126**, e2020JD033681.
- 1641 Wunsch, K. E., and A. C. Didlake, 2018: Analyzing tropical cyclone structures during
1642 secondary eyewall formation using aircraft in situ observations. *Mon. Wea. Rev.*, **146**, 3977-
1643 3993.
- 1644 Xu, J., and Y. Wang, 2010a: Sensitivity of the simulated tropical cyclone inner-core size to
1645 the initial vortex size. *Mon. Wea. Rev.*, **138**, 4135-4157.
- 1646 —, 2010b: Sensitivity of tropical cyclone inner-core size and intensity to the radial distribution
1647 of surface entropy flux. *J. Atmos. Sci.*, **67**, 1831-1852.
- 1648 Yamada, H., and Coauthors, 2021: The double warm-core structure of Typhoon Lan (2017)
1649 as observed through the first Japanese eyewall-penetrating aircraft reconnaissance. *Journal*
1650 *of the Meteorological Society of Japan. Ser. II*, **99**, 1297-1327.

- 1651 Yang, L., X. Cheng, X. Huang, J. Fei, and X. Li, 2020: Effects of air - sea interaction on the
1652 eyewall replacement cycle of Typhoon Sinlaku (2008): Verification of numerical simulation.
1653 *Earth and Space Science*, **7**, e2019EA000763.
- 1654 Yang, X.-w., Y.-q. Wang, H. Wang, J. Xu, and R.-f. Zhan, 2024: Effect of the Initial Vortex
1655 Structure on Intensity Change During Eyewall Replacement Cycle of Tropical Cyclones: A
1656 Numerical Study. *Journal of Tropical Meteorology*, **30**, 106-117.
- 1657 Yang, Y.-T., H.-C. Kuo, E. A. Hendricks, and M. S. Peng, 2013: Structural and intensity
1658 changes of concentric eyewall typhoons in the western North Pacific basin. *Mon. Wea. Rev.*,
1659 **141**, 2632-2648.
- 1660 Yang, Y. T., H. C. Kuo, S. Tsujino, B. F. Chen, and M. S. Peng, 2021: Characteristics of the
1661 Long - Lived Concentric Eyewalls in Tropical Cyclones. *J. Geophys. Res. Atmos.*, **126**,
1662 e2020JD033703.
- 1663 Yu, C.-L., and A. C. Didlake, 2019: Impact of stratiform rainband heating on the tropical
1664 cyclone wind field in idealized simulations. *J. Atmos. Sci.*, **76**, 2443-2462.
- 1665 Yu, C.-L., A. C. Didlake Jr, and F. Zhang, 2022: Updraft Maintenance and Axisymmetrization
1666 during Secondary Eyewall Formation in a Model Simulation of Hurricane Matthew (2016). *J.*
1667 *Atmos. Sci.*, **79**, 1105-1125.
- 1668 Yu, C.-L., A. C. Didlake, F. Zhang, and R. G. Nystrom, 2021a: Asymmetric rainband
1669 processes leading to secondary eyewall formation in a model simulation of Hurricane
1670 Matthew (2016). *J. Atmos. Sci.*, **78**, 29-49.
- 1671 Yu, C. L., A. C. Didlake Jr, J. D. Kepert, and F. Zhang, 2021b: Investigating Axisymmetric
1672 and Asymmetric Signals of Secondary Eyewall Formation Using Observations - Based
1673 Modeling of the Tropical Cyclone Boundary Layer. *J. Geophys. Res. Atmos.*, **126**,
1674 e2020JD034027.
- 1675 Zhang, F., D. Tao, Y. Q. Sun, and J. D. Kepert, 2017a: Dynamics and predictability of
1676 secondary eyewall formation in sheared tropical cyclones. *Journal of Advances in Modeling*
1677 *Earth Systems*, **9**, 89-112.
- 1678 Zhang, J. A., and E. W. Uhlhorn, 2012: Hurricane sea surface inflow angle and an
1679 observation-based parametric model. *Mon. Wea. Rev.*, **140**, 3587-3605.
- 1680 Zhang, J. A., and R. F. Rogers, 2019: Effects of parameterized boundary layer structure on
1681 hurricane rapid intensification in shear. *Mon. Wea. Rev.*, **147**, 853-871.
- 1682 Zhang, J. A., R. F. Rogers, and V. Tallapragada, 2017b: Impact of parameterized boundary
1683 layer structure on tropical cyclone rapid intensification forecasts in HWRF. *Mon. Wea. Rev.*,
1684 **145**, 1413-1426.
- 1685 Zhang, J. A., R. F. Rogers, D. S. Nolan, and F. D. Marks, 2011: On the characteristic height
1686 scales of the hurricane boundary layer. *Mon. Wea. Rev.*, **139**, 2523-2535.
- 1687 Zhang, J. A., D. S. Nolan, R. F. Rogers, and V. Tallapragada, 2015: Evaluating the impact
1688 of improvements in the boundary layer parameterization on hurricane intensity and structure
1689 forecasts in HWRF. *Mon. Wea. Rev.*, **143**, 3136-3155.
- 1690 Zhang, J. A., R. F. Rogers, P. D. Reasor, and J. Gamache, 2023: The mean kinematic
1691 structure of the tropical cyclone boundary layer and its relationship to intensity change. *Mon.*
1692 *Wea. Rev.*, **151**, 63-84.
- 1693 Zhang, J. A., R. F. Rogers, P. D. Reasor, E. W. Uhlhorn, and F. D. Marks, 2013: Asymmetric
1694 hurricane boundary layer structure from dropsonde composites in relation to the
1695 environmental vertical wind shear. *Mon. Wea. Rev.*, **141**, 3968-3984.
- 1696 Zhang, J. A., K. B. Katsaros, P. G. Black, S. Lehner, J. R. French, and W. M. Drennan, 2008:
1697 Effects of roll vortices on turbulent fluxes in the hurricane boundary layer. *Bound.-Layer*
1698 *Meteor.*, **128**, 173-189.
- 1699 Zhang, J. A., and Coauthors, 2018: Evaluating the impact of improvement in the horizontal
1700 diffusion parameterization on hurricane prediction in the operational Hurricane Weather

1701 Research and Forecast (HWRF) Model. *Wea. Fore.*, **33**, 317-329.
1702 Zhu, X.-S., H. Yu, and Y. Wang, 2022: Downwind development in a stationary band complex
1703 leading to the secondary eyewall formation in the simulated Typhoon Soudelor (2015). *Mon.*
1704 *Wea. Rev.*, **150**, 2459-2483.
1705